Theoretical Physics 5 : WS 2024/2025 Exercise sheet 3

04.11.2024

Exercise 0.

How much time did it take to complete this exercise sheet?

Exercise 1. (20 points): Fermion wave functions & operators

- a) (*5 p.*) Explicity write the following *N*-fermion wave functions in terms of singlefermion wave functions $\psi_i(x_j)$:
	- \bullet $\Phi_{0011\cdots0}(x_1,x_2)$
	- \bullet $\Phi_{1101\cdots0}(x_1, x_2, x_3)$
	- $\Phi_{1210\cdots0}(x_1,x_2,x_3,x_4)$
- b) (*15 p.*) The creation and annihilation operators for fermions satisfy

$$
\{c_i, c_j^{\dagger}\} = \delta_{ij}, \qquad \{c_i, c_j\} = \{c_i^{\dagger}, c_j^{\dagger}\} = 0.
$$

Show that for a system of fermions the particle number operator $n = \sum_i c_i^{\dagger}$ $\int_{i}^{1}c_{i}$ commutes with the Hamiltonian

$$
\mathcal{H} = \sum_{i,j} \langle i | H_0 | j \rangle c_i^{\dagger} c_j + \frac{1}{2} \sum_{i,j,k,l} \langle i,j | V | k,l \rangle c_i^{\dagger} c_j^{\dagger} c_l c_k.
$$

Hint: To save some effort, after you show that $\left[n, c \right]$ $\begin{bmatrix} \n\dot{c} \\
\dot{c} \\
\dot{c} \\
\end{bmatrix} = 0$ you can try to manipulate $[n, c_i^{\dagger}]$ $\frac{1}{i}c_j^\dagger$ $\left[\frac{1}{2}c_{l}c_{k}\right]$ into a form whose terms are all similar to $\left[n, c_{i}^{\dagger}\right]$ $\left[\begin{matrix}c_j\end{matrix}\right]$.

Exercise 2. (80 points): High-density electron gas and perturbation theory

The Hamiltonian of a homogeneous electron gas is given by $\mathcal{H} = \mathcal{H}_0 + \mathcal{H}_1$, with

$$
\mathcal{H}_0 = \sum_{\mathbf{k}} \sum_{s} \frac{\hbar^2 k^2}{2m} c_{\mathbf{k},s}^{\dagger} c_{\mathbf{k},s} \quad \text{and} \quad \mathcal{H}_1 = \frac{e^2}{2V} \sum_{\mathbf{k}, \mathbf{p}} \sum_{\mathbf{q} \neq \mathbf{0}} \sum_{s,s'} \frac{4\pi}{q^2} c_{\mathbf{k}+\mathbf{q},s}^{\dagger} c_{\mathbf{p}-\mathbf{q},s'}^{\dagger} c_{\mathbf{p},s'} c_{\mathbf{k},s},
$$

where *m* is the electron mass, $k \equiv |\mathbf{k}|$, and $q \equiv |\mathbf{q}|$.

In the lecture notes you have seen that it is possible to rewrite this Hamiltonian using dimensionless variables, including $r_s \equiv r_0/a_0$, where r_0 is the interparticle spacing. We claimed that in the high-density limit, where $r_s \to 0$, \mathcal{H}_1 is a perturbation to \mathcal{H}_0 .

In this exercise you are going to put this claim to the test by calculating the correction from \mathcal{H}_1 to the ground state energy, and checking whether this correction is indeed small as compared to the ground state energy.

a) (*10 p.*) Derive the magnitude of the Fermi momentum and show that it is given by

$$
k_F = \left(\frac{9\pi}{4}\right)^{1/3} \frac{1}{r_0}.
$$

Here, the interparticle spacing r_0 is defined via $V/N = (4\pi/3)r_0^3$, where *N* is the number of electrons.

b) (*20 p.*) Let us denote the (unperturbed) ground state and the corresponding ground state energy by $|\Psi_0\rangle$ and $E^{(0)} = \langle \Psi_0 | \mathcal{H}_0 | \Psi_0 \rangle$, respectively. Show that

$$
\frac{E^{(0)}}{N} = \frac{3}{5} \frac{\hbar^2 k_F^2}{2m}
$$

Hints:

- You do not need an explicit form for $|\Psi_0\rangle$. Think about what the Fermi momentum means, physically, and how this relates to the sum over *k* of the creation and annihilation operators.
- You might find it useful to use the Heaviside step function,

$$
\Theta(x) = \begin{cases} 1 & \text{if } x > 0, \\ 0 & \text{if } x < 0. \end{cases}
$$

• In the limit of $V \to \infty$ you can replace

$$
\sum_{\mathbf{k},s} f_s(\mathbf{k}) \to \frac{V}{(2\pi)^3} \sum_s \int \mathrm{d}^3 k \, f_s(\mathbf{k}).
$$

c) (*20 p.*) Recall that the first-order correction to the energy is given by

$$
E^{(1)} = \langle \Psi_0 | \mathcal{H}_1 | \Psi_0 \rangle \,. \tag{1}
$$

In order to calculate this correction, it is helpful to consider which combinations of states the operator \mathcal{H}_1 can act on.

Show that

$$
\sum_{k,p} \sum_{q \neq 0} \sum_{s,s'} \langle \Psi_0 | c_{k+q,s}^\dagger c_{p-q,s'}^{\dagger} c_{p,s'} c_{k,s} | \Psi_0 \rangle
$$

=
$$
-2 \left(\frac{V}{(2\pi)^3} \right)^2 \int d^3k \Theta(k_F - k) \int d^3q \Theta(k_F - |\mathbf{k} + \mathbf{q}|),
$$

and then fill in the constants and the $1/q^2$ factor from \mathcal{H}_1 to get an integral expression for $E^{(1)}$.

Hints:

- Consider the order of the annihilation and creation operators acting on $|\Psi_0\rangle$. Since the initial and final state are both the ground state, can you draw any conclusions about the relationship between k , q and p ? What about the relationship between *s* and *s*^{'?} Can you express these restrictions with delta functions?
- The annihilation and creation operators must all act on occupied states, otherwise the matrix element of the transition would be zero. What does it mean for a state with momentum *k* to be occupied? Can you express these restrictions with Heaviside functions?
- Remember the $V \to \infty$ limit from the previous question!

d) (*20 p.*) Perform the integrations to show that

$$
\frac{E^{(1)}}{N} = -\frac{e^2}{4\pi}3k_F.
$$

Hint: If you change integration variables from $k \to P$, using $k = P - q/2$, you can visualize the *P* integral over the two step functions as the volume of the overlap region of two spheres with radius k_F , whose centers are at a distance q from each other:

Remember that these are **spheres** and not circles. **Explain why the** *P* **integral over the** Θ **functions is equivalent to this volume**, and then try to calculate such a volume by integrating in spherical coordinates!

After you get the *P* integration out of the way, the *q* integration should be more straightforward.

e) (*10 p.*) Express $(E^{(0)} + E^{(1)})/N$ in terms of $r_s = r_0/a_0$, where $a_0 = \hbar^2/(me^2)$. Fill in physical values for the constants. Is $E^{(1)}$ indeed small as compared to $E^{(0)}$ in the high-density limit, i.e. is it valid to use perturbation theory here?