# Exercise sheet 11 <br> Theoretical Physics 5 : SS 2023 

26.06.2023

Write your name and your tutor's name on every page you hand in. Please staple said pages together.

## Exercise 0.

How much time did you take to complete this homework sheet?

## Exercise 1. (100 points): <br> Dirac particle in a spherical potential well

Consider the Dirac equation,

$$
\begin{equation*}
\left[\boldsymbol{\alpha} \cdot \boldsymbol{p} c+\beta m_{0} c^{2}\right] \psi(\boldsymbol{r})=[E-V(r)] \psi(\boldsymbol{r}) \tag{1}
\end{equation*}
$$

in a spherical potential well,

$$
V(r)= \begin{cases}-V_{0} & \text { for } \quad r \leq R  \tag{2}\\ 0 & \text { for } \quad r>R\end{cases}
$$

a) (25 p.) Show that

$$
\begin{equation*}
\boldsymbol{\alpha} \cdot \boldsymbol{p}=-i\left(\boldsymbol{\alpha} \cdot \hat{\boldsymbol{e}}_{r}\right)\left(\hbar \frac{\partial}{\partial r}+\frac{\hbar}{r}-\frac{\beta}{r} K\right) \tag{3}
\end{equation*}
$$

with $K=\beta(\boldsymbol{\Sigma} \cdot \boldsymbol{L}+\hbar)$ and $\hat{\boldsymbol{e}}_{r}=\boldsymbol{r} / r$.
Hint: use $\boldsymbol{\nabla}=\hat{\boldsymbol{e}}_{r}\left(\hat{\boldsymbol{e}}_{r} \cdot \boldsymbol{\nabla}\right)-\hat{\boldsymbol{e}}_{r} \times(\hat{\boldsymbol{e}} \times \boldsymbol{\nabla})$.
b) (25 p.) Use the ansatz

$$
\begin{equation*}
\psi(\boldsymbol{r})=\binom{g(r) \phi_{j l_{A} m}(\theta, \phi)}{i f(r) \phi_{j l_{B} m}(\theta, \phi)}, \tag{4}
\end{equation*}
$$

where $\phi_{j l_{A / B} m}$ are the eigenstates of the $\boldsymbol{\sigma} \cdot \boldsymbol{L}$ operator,

$$
\begin{aligned}
& (\boldsymbol{\sigma} \cdot \boldsymbol{L}) \phi_{j l_{A} m}=-\hbar(\kappa+1) \phi_{j l_{A} m}, \\
& (\boldsymbol{\sigma} \cdot \boldsymbol{L}) \phi_{j l_{B} m}=\hbar(\kappa-1) \phi_{j l_{B} m},
\end{aligned}
$$

to find the differential equations for $G(r)$ and $F(r)$. They are related to $g(r)$ and $f(r)$ as

$$
f(r)=\frac{F(r)}{r}, \quad g(r)=\frac{G(r)}{r} .
$$

c) (25 p.) For

$$
k^{2}=\frac{1}{(\hbar c)^{2}}\left[\left(E+V_{0}\right)^{2}-m_{0}^{2} c^{4}\right]>0
$$

the general solution is given by

$$
\begin{aligned}
& G(r)=r\left[a_{1} j_{l_{A}}(k r)+a_{2} y_{l_{A}}(k r)\right] \\
& F(r)=\frac{\kappa}{|\kappa|} \frac{\hbar c k r}{E+V_{0}+m_{0} c^{2}}\left[a_{1} j_{l_{B}}(k r)+a_{2} y_{l_{B}}(k r)\right],
\end{aligned}
$$

where $j_{l}$ and $y_{l}$ are the spherical Bessel functions of the first and second kind. For

$$
\tilde{k}^{2}=-\frac{1}{(\hbar c)^{2}}\left[\left(E+V_{0}\right)^{2}-m_{0}^{2} c^{4}\right]>0
$$

the general solution is given by:

$$
\begin{aligned}
& G(r)=r \sqrt{\frac{2 \tilde{k} r}{\pi}}\left[b_{1} K_{l_{A}+1 / 2}(\tilde{k} r)+b_{2} I_{l_{A}+1 / 2}(\tilde{k} r)\right] \\
& F(r)=\frac{\hbar c \tilde{k} r}{E+V_{0}+m_{0} c^{2}} \sqrt{\frac{2 \tilde{k} r}{\pi}}\left[-b_{1} K_{l_{B}+1 / 2}(\tilde{k} r)+b_{2} I_{l_{B}+1 / 2}(\tilde{k} r)\right]
\end{aligned}
$$

where $K_{l+1 / 2}$ and $I_{l+1 / 2}$ are the modified Bessel functions. Furthermore,

$$
l_{A}=\left\{\begin{array}{lll}
j+\frac{1}{2} & \text { for } \quad & \kappa=j+\frac{1}{2} \\
j-\frac{1}{2} & \text { for } \quad & \kappa=-\left(j+\frac{1}{2}\right)
\end{array}\right.
$$

and

$$
l_{B}=\left\{\begin{array}{lll}
j-\frac{1}{2} & \text { for } \quad & \kappa=j+\frac{1}{2} \\
j+\frac{1}{2} & \text { for } \quad & \kappa=-\left(j+\frac{1}{2}\right)
\end{array}\right.
$$

Determine the transcendental equation that gives the energy spectrum of bound states for which $E-m_{0} c^{2}>-V_{0}$ and $-m_{0} c^{2}<E<m_{0} c^{2}$. You do not need to solve this equation.
Hint: your solution must be normalizable. Consider limits of the special functions (Bessel, etc.) - do they exclude certain solutions?
d) (25 p.) By exploiting the continuity condition at $r=R$, derive the following relation for $s$-states $(l=0)$, which correspond to $j=\frac{1}{2}$ and $\kappa=-1$,

$$
\begin{equation*}
\frac{k R \sin (k R)}{\sin (k R)-k R \cos (k R)}=\frac{k}{\tilde{k}_{0}} \frac{e^{-\tilde{k}_{0} R}}{e^{-\tilde{k}_{0} R}\left(1+\frac{1}{\tilde{k}_{0} R}\right)} \frac{E+m_{0} c^{2}}{E+V_{0}+m_{0} c^{2}}, \tag{5}
\end{equation*}
$$

where

$$
\tilde{k}_{0}=\frac{1}{\hbar c} \sqrt{m_{0}^{2} c^{4}-E^{2}}
$$

If you did not solve part c), take the above equation as your point of departure. Then, rewrite the above equation into

$$
\begin{align*}
1=\tan & \left(\frac{R}{\hbar c} \sqrt{\left(E+V_{0}\right)^{2}-m_{0}^{2} c^{4}}\right) \sqrt{\frac{E+V_{0}+m_{0} c^{2}}{E+V_{0}-m_{0} c^{2}}} \\
& \times\left\{\frac{\hbar c}{R}\left[\frac{1}{E+m_{0} c^{2}}-\frac{1}{E+V_{0}+m_{0} c^{2}}\right]-\sqrt{\frac{m_{0} c^{2}-E}{m_{0} c^{2}+E}}\right\} . \tag{6}
\end{align*}
$$

These equations relate the energy eigenvalues of the $s$-states to the properties of the spherical potential well. It is rather remarkable and nice that we can (still) find analytic solutions!
Hint: If necessary, you can use the following Rayleigh formulas for (spherical) Bessel functions:

$$
\begin{aligned}
& j_{n}(x)=(-1)^{n} x^{n}\left(\frac{1}{x} \frac{\mathrm{~d}}{\mathrm{~d} x}\right)^{n}\left(\frac{\sin (x)}{x}\right), \\
& y_{n}(x)=-(-1)^{n} x^{n}\left(\frac{1}{x} \frac{\mathrm{~d}}{\mathrm{~d} x}\right)^{n}\left(\frac{\cos (x)}{x}\right), \\
& i_{n}(x)=x^{n}\left(\frac{1}{x} \frac{\mathrm{~d}}{\mathrm{~d} x}\right)^{n}\left(\frac{\sinh (x)}{x}\right), \\
& k_{n}(x)=(-1)^{n} x^{n}\left(\frac{1}{x} \frac{\mathrm{~d}}{\mathrm{~d} x}\right)^{n}\left(\frac{e^{-x}}{x}\right),
\end{aligned}
$$

with

$$
\begin{equation*}
i_{n}(x)=\sqrt{\frac{\pi}{2 x}} I_{n+1 / 2}(x) \quad \text { and } \quad k_{n}(x)=\sqrt{\frac{\pi}{2 x}} K_{n+1 / 2}(x) . \tag{7}
\end{equation*}
$$

