Exercise sheet 3

Theoretical Physics 3: QM WS2022/2023

Lecturer: Prof. Dr. M. Vanderhaeghen

09.11.2022

Exercise 1. Harmonic oscillator: ladder operators (35 points)

Consider a quantum harmonic oscillator, the time-independent ground state wave function of which is given by:

$$\psi_0(x) = \sqrt[4]{\frac{m\omega}{\pi\hbar}} e^{-\frac{m\omega x^2}{2\hbar}} \equiv \alpha e^{-\frac{y^2}{2}}$$

Where, for further simplicity, we have introduced $\alpha = \left(\frac{m\omega}{\pi\hbar}\right)^{\frac{1}{4}}$ and dimensionless variable $y = \sqrt{\frac{m\omega}{\hbar}}x$.

a) (5 p.) Using explicit definition of the raising ladder operator:

$$\hat{a} = \frac{1}{\sqrt{2\hbar\omega m}}(-i\hat{p} + m\omega\hat{x}) \equiv \frac{1}{\sqrt{2}}\left(-\frac{\mathrm{d}}{\mathrm{d}y} + y\right)$$

Derive expression for the first excited state wave function ψ_1 and check its orthogonality to ψ_0 .

- b) (20 p.) Compute $\langle x \rangle, \langle p \rangle, \langle x^2 \rangle$ and $\langle p^2 \rangle$, for the states ψ_0 and ψ_1 by explicit integration.
- c) (5 p.) Check the uncertainty principle for these two states.
- d) (5 p.) Compute expectation values of the kinetic energy $\langle T \rangle$ and the potential energy $\langle V \rangle$. Check these to sum up to $\langle H \rangle$.

Exercise 2. Harmonic oscillator: power series method (65 points)

The quantum harmonic oscillator problem can be solved using the power series method. One starts with the stationary Schrödinger equation $(\psi'' \equiv \frac{d^2\psi}{dx^2})$:

$$-\frac{\hbar^2}{2m} \psi''(x) + \frac{1}{2} m\omega^2 x^2 \psi(x) = E \psi(x)$$

a) (10 p.) To simplify the initial problem, rewrite the equation using the dimensionless quantities

$$y = \sqrt{\frac{m\omega}{\hbar}}x, \quad \varepsilon = E/\hbar\omega$$

Further on, define $\varphi(y) = c\psi(x)$ and find c, such that $\varphi(y)$ is normalized.

b) (10 p.) Investigate the asymptotical behaviour of the equation for large y. Show that for $y \to \infty$:

$$\varphi(y) \sim e^{-\frac{y^2}{2}}$$

c) (10 p.) We can now explicitly isolate the asymptotic behaviour of the unknown function:

$$\varphi(y) = h(y) e^{-\frac{y^2}{2}}$$

Derive the following equation on h(y):

$$h'' - 2yh' + (2\varepsilon - 1)h = 0$$

d) (15 p.) At this point, assume that h(y) can be written as an infinite power series in y:

$$h(y) = \sum_{m=0}^{\infty} a_m y^m$$

Derive the recurrence relation between the coefficients a_m and show that there are two sets of independent solutions (even and odd).

- e) (15 p.) Prove that, in order for the wave function to be finite and normalizable, one has to imply that the infinite series must be "cut off" at some finite integer n: $a_{m>n} = 0$. (*Hint*: consider Maclaurin expansion of e^{y^2} and compare it to the series behaviour for large y).
- f) (5 p.) Using the previous conclusion, show that the energy is quantized $E_n = (n + \frac{1}{2})\hbar\omega$.

The obtained polynomials $h_n(y)$ are proportional to Hermite polynomials $H_n(y)$. The orthonormal set of solutions of the initial stationary Schrödinger equation then reads:

$$\psi_n(x) = \left(2^n n! \sqrt{\frac{\pi \hbar}{m\omega}}\right)^{-\frac{1}{2}} H_n\left(\sqrt{\frac{m\omega}{\hbar}}x\right) e^{-\frac{m\omega x^2}{2\hbar}}, \quad n = 0, 1, 2, \dots$$

(Bonus) Exercise 3. Supersymmetric QM (50 points)

We consider in this exercise a generalization of the raising and lowering operator method. For a given potential $V_{-}(x)$, the idea is to construct a partner (also called supersymmetric) potential $V_{+}(x)$ which has the same energy eigenvalues, except for the ground state. Without loss of generality, we can shift the potential $V_{-}(x)$ so that the corresponding ground state $\psi_{0}(x)$ has zero energy $E_{0}^{-}=0$.

a) (10 p.) Show that the Hamiltonian $\hat{H}_{-} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V_{-}$ can be written in the form:

$$\hat{H} = \hat{A}^{\dagger} \hat{A}$$

With:

$$\hat{A}^{+} \equiv \frac{\hbar}{\sqrt{2m}} \left(-\frac{d}{dx} - \frac{\psi'_{0}}{\psi_{0}} \right)$$

$$\hat{A} \equiv \frac{\hbar}{\sqrt{2m}} \left(\frac{d}{dx} - \frac{\psi'_{0}}{\psi_{0}} \right)$$

And $\psi_0' \equiv \frac{d}{dx}\psi_0$.

b) (10 p.) Consider the partner Hamiltonian operator $\hat{H}_{+} = \hat{A}\hat{A}^{+}$ which can also be defined as $\hat{H}_{+} = -\frac{\hbar^{2}}{2m}\frac{d^{2}}{dx^{2}} + V_{+}$. Show that the partner potential V_{+} is related to V_{-} as follows:

$$V_{+} = V_{-} - \frac{\hbar^2}{m} \frac{d}{dx} \left(\frac{\psi_0'}{\psi_0} \right)$$

c) (15 p.) Show that \hat{H}_{-} and \hat{H}_{+} have the same spectrum, except for the ground state (*Hint*: Consider the states $\hat{A}\psi_{n}^{-}$ and $\hat{A}^{+}\psi_{n}^{+}$ with ψ_{n}^{\pm} being eigenstates of \hat{H}_{\pm}). Write the eigenstates ψ_{n}^{+} and energies E_{n}^{+} in terms of ψ_{n}^{-} and E_{n}^{-} .

d) (15 p.) Consider a particle in the infinite square potential well:

$$V_{-}(x) = \begin{cases} V_{0} & \text{for } 0 \le x \le a \\ +\infty & \text{otherwise,} \end{cases}$$

Find V_0 such that the ground state has zero energy. Derive the partner potential $V_+(x)$. Write down the properly normalized eigenstates $\psi_n^+(x)$. Explain why the existence of partner potentials may be useful.