

Exercise sheet 6  
Theoretical Physics 5 : WS 2021/2022  
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**Exercise 0.**

How much time did it take to complete the task?

**Exercise 1. (50 points) : Spontaneous symmetry breaking**

Consider the following Lagrangian for a complex scalar field  $\phi = (\phi_1 + i\phi_2)/\sqrt{2}$

$$\mathcal{L} = (\partial^\mu \phi^*)(\partial_\mu \phi) + \mu^2 \phi^* \phi - \frac{\lambda}{2} (\phi^* \phi)^2,$$

Clearly, for  $\mu^2 = -m_0^2 c^2 / \hbar^2$  and  $\lambda = 0$ , one recovers the familiar Klein-Gordon Lagrangian. Moreover, this Lagrangian is also invariant under the phase transformation  $\phi \rightarrow e^{i\alpha} \phi$ .

- a) (10 p.) Treating  $\phi$  and  $\phi^*$  as independent, show that the corresponding Hamiltonian is given by

$$H = \int d^3x \left( \frac{1}{c^2} \dot{\phi}^* \dot{\phi} + \vec{\nabla} \phi^* \cdot \vec{\nabla} \phi - \mu^2 \phi^* \phi + \frac{\lambda}{2} (\phi^* \phi)^2 \right).$$

*Hint:* Remember that the conjugate momentum  $\pi$  of a field  $\phi$  is given by  $\pi = \partial \mathcal{L} / \partial \dot{\phi}$  and the Hamiltonian by  $H = \int d^3x (\pi \dot{\phi} + \pi^* \dot{\phi}^* - \mathcal{L})$ .

- b) (20 p.) Find the condition on classical field configurations  $\phi_0(x)$  which assures minimization of the energy. Show that there is an infinite set of such configurations and that they are all related by the phase transformation. *Hint:* Try first to minimize the “kinetic” term  $\frac{1}{2} \dot{\phi}^* \dot{\phi} + \vec{\nabla} \phi^* \cdot \vec{\nabla} \phi$ , then minimize the “potential” term  $-\mu^2 \phi^* \phi + \frac{\lambda}{2} (\phi^* \phi)^2$ .
- c) (30 p.) Suppose that the system is near the minimum  $\phi_0 = \mu / \sqrt{\lambda}$ . Then it is convenient to define

$$\phi(x) = \phi_0 + \frac{1}{\sqrt{2}} [\sigma(x) + i\theta(x)],$$

where the fields  $\sigma(x)$  and  $\theta(x)$  describe (real) fluctuations around the minimum. Rewrite  $\mathcal{L}$  in terms of these fluctuations. What are the masses of the  $\sigma$  and  $\theta$  fields? *Hint*: Remember that the mass can be read off the coefficient quadratic in the field, *e.g.* for a real scalar field  $\Phi$  with mass  $m_0$  we had the quadratic term  $-\frac{1}{2} \frac{m_0^2 c^2}{\hbar^2} \Phi^2$ .

## Exercise 2. (50 points) : Relativistic electrons in a constant magnetic field

To study the behavior of electrons in a constant magnetic field, we have to solve the stationary-state Dirac equation

$$\left(-i\hbar c \vec{\alpha} \cdot \vec{D} + \beta m_0 c^2\right) \psi = E\psi$$

where we have used minimal substitution  $\vec{\nabla} \rightarrow \vec{D} \equiv \vec{\nabla} - ie\vec{A}/\hbar c$ .

a) (20 p.) Verify that

$$\left(\vec{\alpha} \cdot \vec{D}\right)^2 = \vec{D}^2 1 + \frac{e}{\hbar c} \vec{\Sigma} \cdot \vec{B}, \quad \text{where} \quad \vec{\Sigma} = \begin{pmatrix} \vec{\sigma} & 0 \\ 0 & \vec{\sigma} \end{pmatrix}.$$

b) (30 p.) For the particular case  $\vec{A} = (0, xB, 0)$  and by considering solutions of the form  $\psi = e^{i(p_y y + p_z z)/\hbar} u(x)$ , show that the energy eigenvalues  $E$  of a relativistic electron in constant magnetic induction  $\vec{B}$  are given by

$$E^2 = m_0^2 c^4 + p_z^2 c^2 + (2n + 1)|eB|\hbar c \pm eB\hbar c, \quad n \in N.$$

*Hint*: Remember that the eigenvalues of the harmonic oscillator operator  $-\partial_x^2 + \omega^2 x^2$  are  $(2n + 1)|\omega|$  with  $n \in N$ .