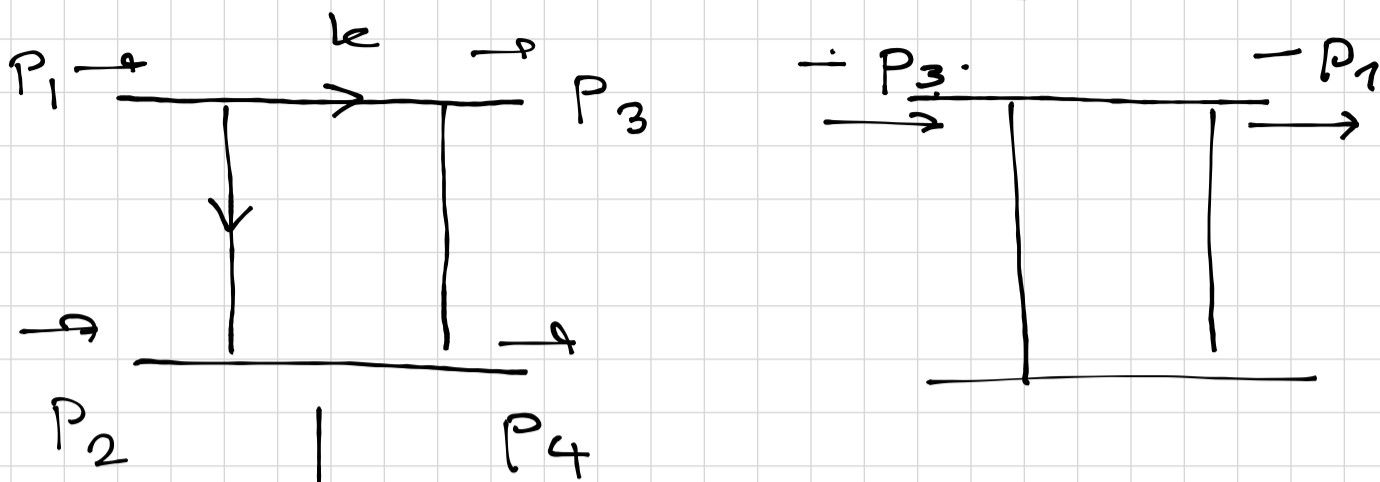


Lecture 11 7/12/20



$$k = (\alpha, \beta, k_{\perp}) = \alpha \cdot n^{\mu} + \beta \bar{n}^{\mu} + (0, \vec{k}_{\perp}, 0)$$

$$p_1^{\mu} \approx (1, \gamma, 0_{\perp}) \quad n, \bar{n} = \frac{\sqrt{s}}{2} (1, 0, 0, \pm 1)$$

$$p_2^{\mu} \approx (\gamma, 1, 0_{\perp}) \quad n^2 = \bar{n}^2 = 0$$

$$\gamma = \frac{m^2}{s} \quad 2n\bar{n} = s \Rightarrow k_{\perp}^2, m^2$$

Method of regions \rightarrow determine the positions of poles on complex d -plane (keeping track of signs of $i\epsilon$)

$\hookrightarrow \int d^2 k_{\perp}$ and $\int d\beta$ splits

$\hookrightarrow \sim \frac{\pi}{m^2}$ \hookrightarrow HE asymptotics

At 1-loop level

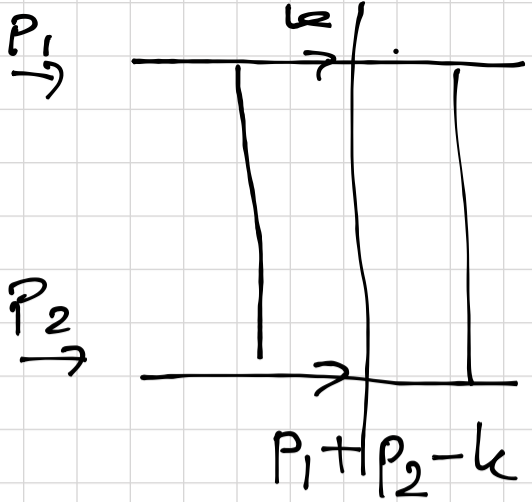
$$A_s^{(4)} \approx -\frac{g^2}{s} \cdot \frac{g^2}{16\pi^2 m^2} \ln\left(-\frac{s}{m^2} - i\epsilon\right)$$

$$A_s^{(2)} \approx -\frac{g^2}{s} \ln\left(\frac{s}{m^2}\right) - i\pi \quad s > m^2$$

\hookrightarrow Was obtained also from a DR:

1). Calculate the \ln part

\rightarrow follows the Cutkosky's rules



$$\rightarrow \frac{1}{k^2 - m^2 + i\epsilon} \frac{1}{(P_1 + P_2 - k)^2 - k^2 + i\epsilon}$$

$$\rightarrow (-2\pi i) \delta(k^2 - m^2) \Theta(k^0)$$

$$(-2\pi i) \delta((P_1 + P_2 - k)^2 - k^2) \Theta(P_1^0 + P_2^0 - k^0)$$

$$\rightarrow i A_s^{(4)}(s + i\epsilon) - i A_s^{(4)}(s - i\epsilon)$$

$$= -2 \operatorname{Im} A_s^{(4)}(s)$$

$$\hookrightarrow \operatorname{Im} A_s^{(4)}(s) = \frac{g^2}{s} \frac{g^2}{16\pi^2 m^2}$$

$$\hookrightarrow \text{DR} \quad \operatorname{Re} A_s^{(4)}(s) = \frac{1}{\pi} \mathcal{P} \int_{4m^2}^{\infty} \frac{ds'}{s' - s} \operatorname{Im} A_s^{(4)}(s')$$

$$\mathcal{P} \int_{4m^2}^{\infty} \frac{ds'}{s' - s} \cdot \frac{1}{s'} = -\frac{1}{s} \operatorname{Ln}\left(\frac{s}{m^2}\right)$$

$$A_s^{(4)}(s) = \frac{1}{\pi} \int \frac{ds'}{s' - s - i\epsilon} \operatorname{Im} A_s^{(4)}(s')$$

$$A_s^{(n)} = -\frac{g^2}{s} \cdot \frac{1}{n!} \left[\frac{g^2}{16\pi^2 m^2} \operatorname{Ln}\left(-\frac{s}{m^2} - i\epsilon\right) \right]^n$$

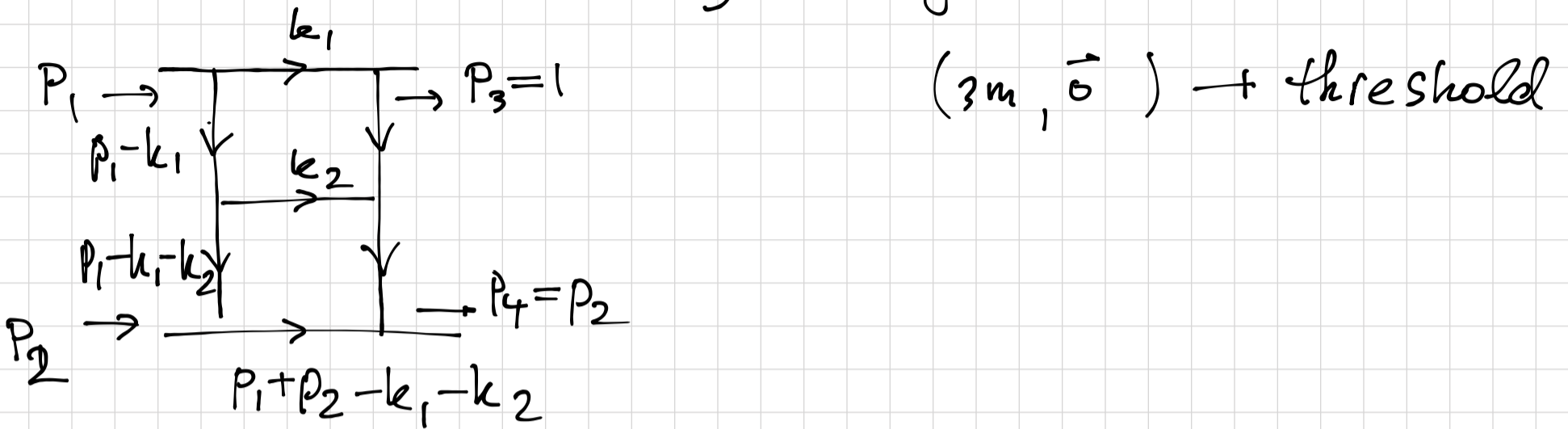
$$\sum_{n=0}^{\infty} A_s^{(n)} = -\frac{g^2}{s} \exp \left[\frac{g^2}{16\pi^2 m^2} \operatorname{Ln}\left(-\frac{s}{m^2} - i\epsilon\right) \right]$$

$$= -\frac{g^2}{s} \cdot \left(-\frac{s}{m^2} - i\epsilon\right)^{\beta}$$

$$\beta = \frac{g^2}{16\pi^2 m^2}$$

Generalize our result to higher order

Let's obtain $A_s^{(6)} \sim g^6$



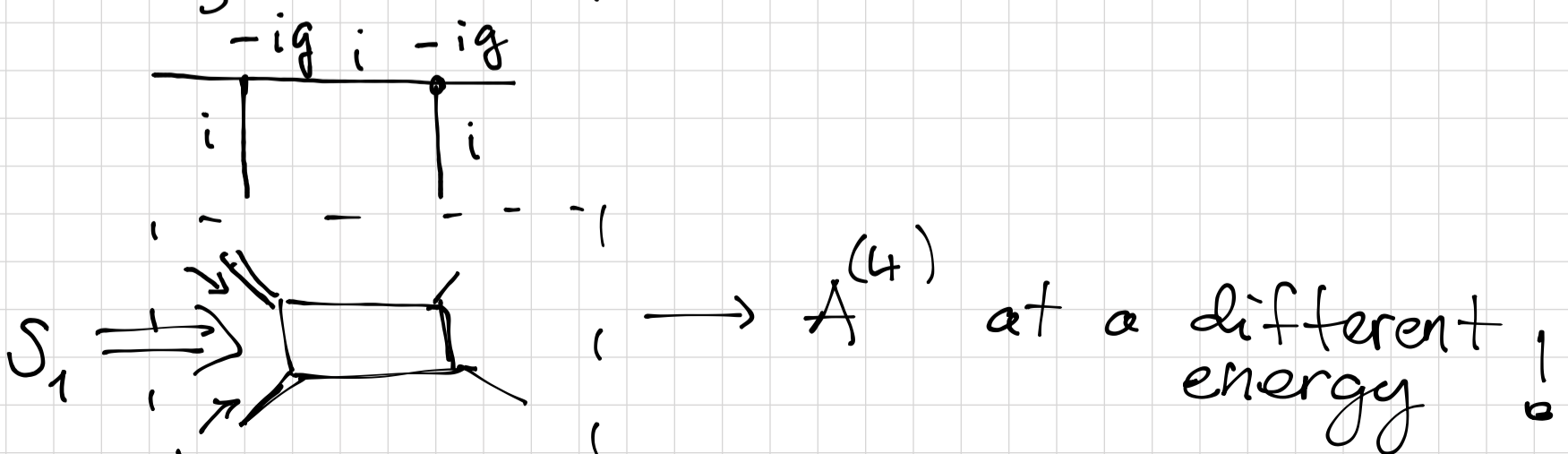
$$= \int \frac{d^4 k_1}{(2\pi)^4} \int \frac{d^4 k_2}{(2\pi)^4} \frac{1}{k_1^2 - m^2 + i\epsilon} \frac{1}{k_2^2 - m^2 + i\epsilon} \frac{1}{(p_1 + p_2 - k_1 - k_2)^2 - m^2 + i\epsilon}$$

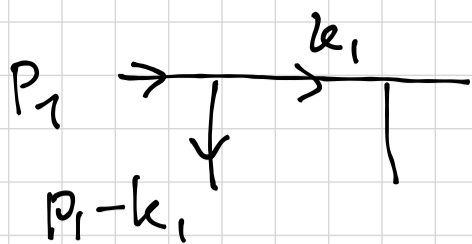
$$\frac{1}{[(p_1 - k_1)^2 - m^2 + i\epsilon]^2}$$

$$\frac{1}{[(p_1 - k_1 - k_2)^2 - m^2 + i\epsilon]^2}$$

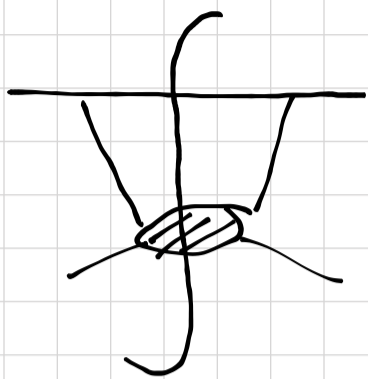
Here we only look asymptotics $s \rightarrow \infty$
 We will use the cutting rules again
 and we use the previously obtained

$A_s^{(4)}$ as input





$$s_1 = (p_1 + p_2 - k_1)^2$$



$$-2 \operatorname{Im} A_S^{(6)}(s) = (-ig)^2 \int \frac{d^4 k_1}{(2\pi)^4} i^3 \times (-2\pi i) \delta(k_1^2 - m^2) \Theta(k_1^0)$$

$$-2 \operatorname{Im} A_S^{(4)}(s_1) \leftarrow \left[i A_S^{(4)}(s_1 + i\epsilon) - i A_S^{(4)}(s_1 - i\epsilon) \right] \left[\frac{1}{(p_1 - k_1)^2 - m^2 + i\epsilon} \right]^2$$

↓

$$\operatorname{Im} A_S^{(6)}(s) = \frac{g^6}{(2\pi)^3 16\pi m^2} \int \frac{d^3 k_1^0 k_1^i d\Omega}{s_1 2k_1} \frac{\delta(k_1^0 - \sqrt{k_1^2 - m^2})}{[(p_1 - k_1)^2 - m^2 + i\epsilon]^2}$$

C.M. $(p_1 + p_2)^\mu = (\sqrt{s}, \vec{0})$ $\Theta(k_1^0) \Theta(s_1 - 4m^2)$

$$s_1 = (p_1 + p_2 - k_1)^2 = s + k_1^2 - 2\sqrt{s} k_1^0$$

$$k_1^0 = \frac{s - s_1 + m^2}{2\sqrt{s}} \xrightarrow{\delta} k_1 = \sqrt{\left(\frac{s - s_1 + m^2}{2\sqrt{s}}\right)^2 - k_1^2} = \frac{\sqrt{[s - (\sqrt{s_1} + m)]^2 [s - (\sqrt{s_1} - m)]^2}}{2\sqrt{s}}$$

$$\hookrightarrow \frac{\Theta(\sqrt{s} - \sqrt{s_1} - m)}{(\sqrt{s} - m)^2} \left. \begin{array}{l} \\ dk_1^0 = \frac{-ds_1}{2\sqrt{s}} \end{array} \right\}$$

$$\Rightarrow \frac{g^6}{(2\pi)^3 16\pi m^2} \frac{1}{2} \int \frac{ds_1}{s_1} \int d\Omega \frac{k_1}{2\sqrt{s}} \frac{1}{[(p_1 - k_1)^2 - m^2 + i\epsilon]^2}$$

$$\equiv \quad \equiv \quad 4m^2$$

$$1. (p_1 - k_1)^2 - m^2 = m^2 - 2p_1 k_1 \approx -\frac{s-s_1}{2} (1 - \cos\theta)$$

$$\begin{aligned} s &>> m^2 \\ s_1 &>> m^2 \end{aligned}$$

$$\frac{k_1}{2\sqrt{s}} \approx \frac{s-s_1}{4s}$$

$$p_1 = (1, 0, 0, 1) \frac{\sqrt{s}}{2}$$

$$k_1 = \frac{s-s_1}{2\sqrt{s}} (1, \hat{k}(\theta, \varphi))$$

$$\Downarrow \frac{g^6}{(2\pi)^3 16\pi m^2} \frac{1}{2} \frac{1}{4s} \int_{4m^2}^{(\sqrt{s}-m)^2} \frac{ds_1}{s_1} (s-s_1) \int \frac{d\varphi d\cos\theta}{\left[\frac{s-s_1}{2}(1-\cos\theta)\right]^2}$$

$$= \frac{g^6}{16\pi^2 16\pi m^2} \frac{1}{m^2} \frac{1}{s} \int_{4m^2}^{(\sqrt{s}-m)^2} \frac{ds_1}{s_1} \quad 2\pi \frac{2}{s-s_1} \left(\frac{1}{m^2} - \frac{1}{s-s_1} \right)$$

$$= \frac{g^2}{s} \cdot \frac{g^2}{16\pi m^2} \cdot \frac{g^2}{16\pi^2 m^2} \cdot \text{Re} \frac{s}{m^2} + \text{corr.}$$

$$\text{Re} A_s^{(6)} = \frac{1}{\pi} \mathcal{P} \int_{gm^2}^{\infty} \frac{ds'}{s'-s} \text{Im} A_s^{(6)}(s')$$

$$= g^2 \left(\frac{g}{4\pi m} \right)^4 \mathcal{P} \int \frac{ds'}{s'(s'-s)} \text{Im} \left(\frac{s'}{m^2} \right)$$

$$\frac{1}{s} \mathcal{P} \int_{gm^2}^{\infty} \left(\frac{ds'}{s'-s} - \frac{ds'}{s'} \right) \text{Im} \frac{s'}{m^2}$$

Sub. $x = \frac{gm^2}{s'} \rightarrow \mathcal{P} \int \dots \rightarrow -2 \int_0^1 \frac{dx}{x-y} \ln \frac{1}{x}$
 $y = \frac{sm^2}{s}$

\downarrow
 $\text{Re } A_s^{(6)} = -\frac{g^2}{s} \left(\frac{g}{4\pi M} \right)^4 \text{Li}_2 \left(\frac{s}{m^2} + i\epsilon \right) + \dots$

Dilogarithm function

$$\text{Li}_2(y) = -\int_0^1 \frac{dt}{t} \ln(1-yt)$$

Properties of Li_2 :

1. $\text{Li}_2(0) = 0$

2. Real for $-\infty < y < 1$; branch cut for $y > 1$

3. $\text{Li}_2(1) = \frac{\pi^2}{6}$

Identities : $\left\{ \begin{array}{l} \text{Li}_2\left(\frac{1}{x}\right) + \text{Li}_2(x) = -\frac{\pi^2}{6} - \frac{1}{2} \ln^2(-x) \\ \text{Li}_2(1-x) + \text{Li}_2(x) = \frac{\pi^2}{6} - \ln x \ln(1-x) \end{array} \right.$

$\stackrel{x}{\parallel}$
 $\text{Li}_2\left(\frac{s}{m^2} + i\epsilon\right) = \text{Li}_2\left(\frac{m^2}{s}\right) - \frac{\pi^2}{6} - \frac{1}{2} \ln^2\left(-\frac{s}{m^2} - i\epsilon\right)$
 $\sim \frac{m^2}{s} \rightarrow 0$

\Downarrow

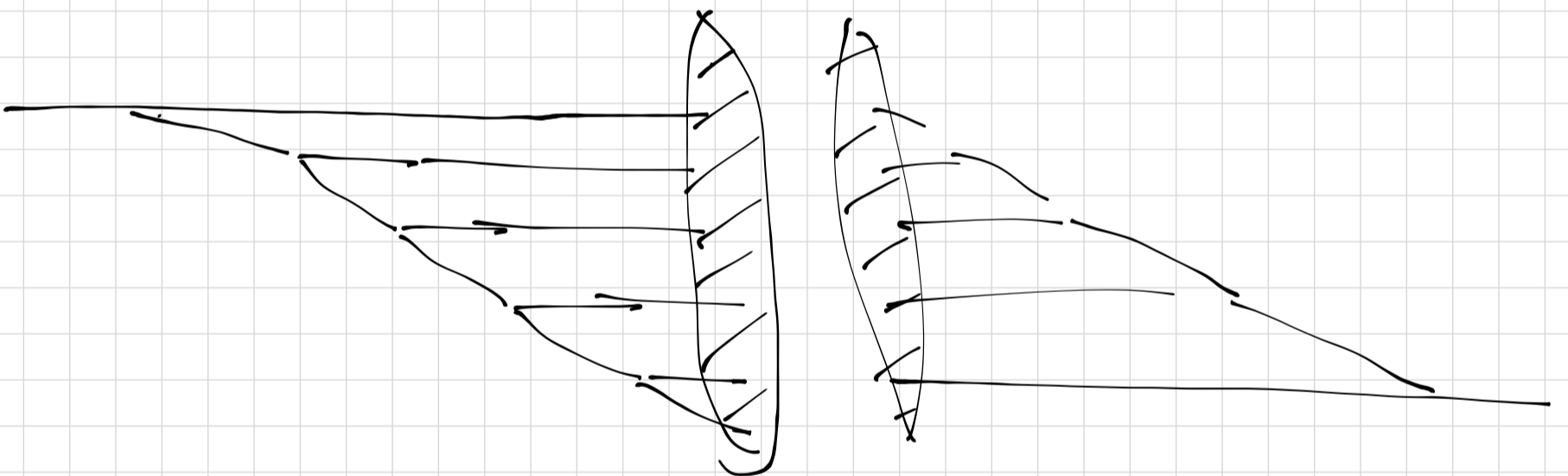
$$A_s^{(6)}(s) \approx -\frac{g^2}{s} \frac{1}{2!} \left(\frac{g}{4\pi M} \right)^4 \ln^2\left(-\frac{s}{m^2} - i\epsilon\right)$$

$$A_s^{(2n)} \approx -\frac{g^2}{s} \cdot \frac{1}{2(n-1)!} \left(\frac{g}{4\pi m}\right)^{2n} \ln^{2(n-1)} \ln\left(-\frac{s}{\mu^2} - i\epsilon\right)$$

$$\text{Im } A^{(n+2)} = g^2 \int \frac{d^4 k_1}{(2\pi)^4} \frac{\delta(k_1^2 - \mu^2) \Theta(\dots)}{\left[(p_1 - k_1)^2 - \mu^2 + i\epsilon\right]^2}$$

• $\text{Im } A^{(n)}(s_1)$

At high energies the interaction between the cloud (shower) of virtual particles



Unitarity of the S-matrix

$$S \approx e^{-iT}$$

$$S_{fi} = \langle f | U(t_2 \rightarrow +\infty, t_1 \rightarrow -\infty) | i \rangle$$

$$S^\dagger S = \mathbb{1}$$

$$S_{fi} \rightarrow (S - \mathbb{1})_{fi}$$

$$S - \mathbb{1} = iT$$

$$S^\dagger S = \mathbb{1} = (\mathbb{1} - iT^\dagger)(\mathbb{1} + iT) = \mathbb{1}$$

$$\rightarrow i(T^\dagger - T) = T^\dagger T$$

$$\langle f | i(T^\dagger - T) | i \rangle = i(2\pi)^4 \delta^4(p_f - p_i)$$

$$T_{fi} = (2\pi)^4 \delta^4(p_f - p_i) A_{fi} \quad \times [A_{f \rightarrow i}^* - A_{i \rightarrow f}]$$

$$\langle f | T^\dagger | i \rangle = \langle i | T | f \rangle^*$$

$$\langle f | T^\dagger T | i \rangle = \sum_x \int \prod_{j \in X} \frac{d^3 p_j}{(2\pi)^3 2E_j} \langle f | T^\dagger | X \rangle \langle X | T | i \rangle$$

$$\prod = \sum_x \int d\pi_x |x\rangle \langle x|$$



$$A_{i \rightarrow f} - A_{f \rightarrow i}^* = i \sum_x \int d\pi_x (2\pi)^4 \delta^4(p_i - p_x) A_{i \rightarrow x} A_{f \rightarrow x}^*$$

Important case: $i=f$ forward scattering
2 → 2

$$A_{i \rightarrow i} - A_{i \rightarrow i}^* = 2iA$$

$$\text{r.h.s.} \rightarrow i \sum_x \int d\pi_x (2\pi)^4 \delta^4(p_i - p_x) |A_{i \rightarrow x}|^2$$

$$\equiv 4E_{cm} |\vec{p}| \sigma_{tot}$$

Optical theorem

$$\text{Im } A_{i \rightarrow i} = 2E_{cm} |\vec{p}| \sigma_{tot}$$

$|i\rangle \rightarrow$ 1-particle state



$$\text{Im } A_{i \rightarrow i} = M_i \cdot \Gamma_{\text{tot}} \sim |\text{diagram}|^2$$

$$\text{Im} \left[\text{diagram} \right]$$

$$\sim |\text{diagram}|^2 \sim \Gamma_{\text{tot}}$$

Resonance $M_R \rightarrow M_R - \frac{i\Gamma_R}{2}$

$$\text{Im}(p^2 - M_R^2) \rightarrow \text{Im}\left(p^2 - \left(M - \frac{i\Gamma}{2}\right)^2\right) = M\Gamma$$

$$\text{Im } A \sim \mathcal{Q} E_{\text{cm}} p \sigma_{\text{tot}}$$

$$A = \sum_{n=0}^{\infty} c_n \frac{1}{n!} g^n$$

$c_n = \text{complex fn.}$

$$\text{Im } A^{(n)} = \frac{1}{n!} g^n c_n = \mathcal{Q} E_{\text{cm}} p \sigma^{(n)}$$

$$A^{(6)} \rightarrow g^6 \left(\text{diagram} \right)^2$$

Loop amplitudes

can be reconstructed from tree-level amplitudes

$$A^{(6)} \sim \int d\pi_x A^{(2)} \cdot A^{(4)}$$

Connects lower order to higher orders

→ we used it for ladder resummation