

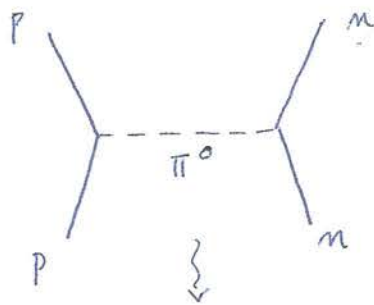
# V.

## INTERACTIONS & ELEMENTS OF FIELD THEORY

INT 1

### \* QUANTUM EXCHANGE: YUKAWA THEORY

YUKAWA (1935) : SHORT RANGE FORCE BETWEEN PROTONS & NEUTRONS



DUE TO EXCHANGE OF A PION

↓  
MASS  $m \approx 135 \text{ MeV}$   
( $\pi^0$ )

EXCHANGED PARTICLE: VIRTUAL QUANTUM  $\Delta E \Delta t \sim \hbar$

↓  
SPIN 0: DESCRIBED BY FIELD  $\Phi(t, \vec{x})$   
SATISFIES KLEIN-GORDON EQ.

$$\left( \frac{\partial^2}{\partial t^2} - \nabla^2 + m^2 \right) \Phi = 0$$

#### SOLUTIONS

- FOR FREE PROPAGATING PARTICLE  $\hat{p}^\mu(E, \vec{p})$   
 $-i(Et - \vec{p} \cdot \vec{x})$

PLANE WAVE  $\Phi(t, \vec{x}) = A e$

$$E^2 = \vec{p}^2 + m^2$$

- FOR STATIC POTENTIAL (P OR n ACT AS SOURCE & SINK)

$$\Phi \rightarrow U(\vec{x})$$

$$\left( -\nabla^2 + m^2 \right) U(\vec{x}) = 0$$

ASSUME SPHERICAL SYMMETRIC POTENTIAL

$$U(\vec{x}) = U(r) \quad r = |\vec{x}|$$

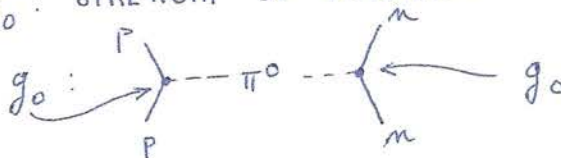
$$\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial U}{\partial r} \right) = m^2 U$$

SOLUTION

$$U(r) = \frac{g_0}{4\pi} \frac{1}{r} e^{-mr}$$

YUKAWA POTENTIAL ('SCREENED'  $\frac{1}{r}$  POTENTIAL)

$\Rightarrow g_0$ : STRENGTH OF POTENTIAL



$\Rightarrow R$ : RANGE OF POTENTIAL  $\sim \frac{1}{m}$

$$R = \frac{(0.197 \text{ GeV fm})}{m} = \frac{0.197}{0.135} \text{ fm}$$

$$= 1.5 \text{ fm}$$

SHORT RANGE  $\Leftrightarrow$  EXCHANGE OF MASSIVE PARTICLE  
OF STRONG INTERACTION

PION WAS OBSERVED IN 1947

\*

## SPIN-0 PROPAGATOR

STATIC POTENTIAL  $U(r)$

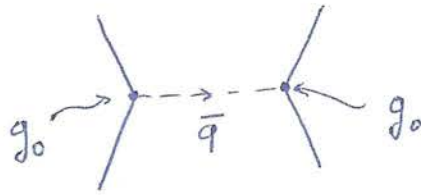
$g_0$ : COUPLING OF PARTICLE TO POTENTIAL

↓

FOURIER TF (TO MOMENTUM SPACE)

$$f(\bar{q}) = g_0 \int d^3 \bar{x} U(r) e^{i\bar{q} \cdot \bar{x}}$$

↳ DESCRIBES EXCHANGE OF VIRTUAL PARTICLE WITH MOMENTUM  $\bar{q}$



$$f(\bar{q}) = \frac{g_0^2}{4\pi} \int_0^\infty dr r^2 \frac{e^{-mr}}{r} \cdot 2\pi \int_0^\pi d\cos\theta e^{i|\bar{q}|r\cos\theta}$$

$$\frac{2}{|\bar{q}|r} \sin(|\bar{q}|r)$$

$$= g_0^2 \int_0^\infty dr \frac{1}{|\bar{q}|} \sin(|\bar{q}|r) e^{-mr}$$

$$= g_0^2 \frac{1}{|\bar{q}|} \frac{1}{2i} \int_0^\infty dr e^{-mr} \left( e^{i|\bar{q}|r} - e^{-i|\bar{q}|r} \right)$$

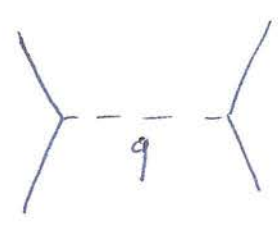
$$= g_0^2 \frac{1}{2i|\bar{q}|} \left\{ -\frac{1}{i|\bar{q}| - m} - \frac{1}{i|\bar{q}| + m} \right\}$$

$$= g_0^2 \frac{1}{|\bar{q}|^2 + m^2}$$

- FOR STATIC POTENTIAL (RANGE  $\sim \frac{1}{m}$ )

$$f(\vec{q}) = \frac{g_0^2}{|\vec{q}|^2 + m^2}$$

- FOR TIME-DEPENDENT EXCHANGE PROCESS



$$q^\mu (q_0, \vec{q})$$

REPLACE  $\vec{q}^2 \rightarrow -q^2 = -q_0^2 + \vec{q}^2$

|| AMPLITUDE :  $\frac{-g_0^2}{q^2 - m^2}$

$\frac{1}{q^2 - m^2} \sim$  SPIN 0 PROPAGATOR

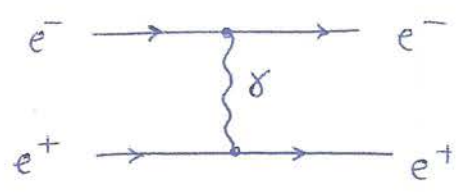
FROM AMPLITUDE FOR A PROCESS

WE CAN CALCULATE CROSS SECTIONS, DECAY RATES, ...

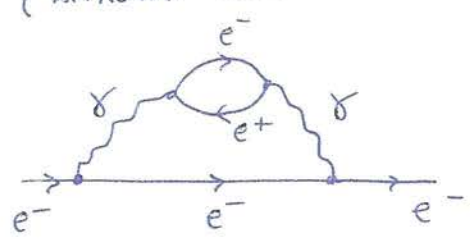
\* FEYNMAN DIAGRAM

AMPLITUDE FOR PROCESS IS GRAPHICALLY DISPLAYED BY A FEYNMAN DIAGRAM (ARROWS INDICATE TIME SENSE)

e.g.

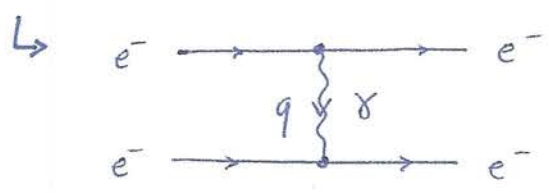


$\sim e^2$



$\sim e^4$

\* ELECTROMAGNETIC INTERACTIONS



$q$ : 4-MOMENTUM OF  $\gamma$

• STRENGTH  $\sim \alpha = \frac{e^2}{4\pi} = \frac{1}{137.036\dots}$   
 $\Downarrow$   
 FINE STRUCTURE CONSTANT

- EXCHANGED PARTICLE : PHOTON (MASSLESS, SPIN 1)  
 $\Downarrow$
- POTENTIAL :  $\infty$  RANGE
  - PROPAGATOR  $\sim \frac{1}{q^2}$

• CROSS SECTION ( $\sim$  SCATTERING PROBABILITY)

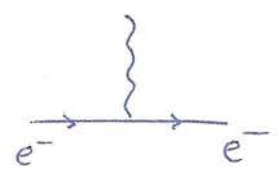
$d\sigma \sim \left| \frac{\alpha}{q^2} \right|^2 = \frac{\alpha^2}{q^4}$  (RUTHERFORD CROSS SECTION)

$\hookrightarrow$  EM COUPLING SMALL  $\Rightarrow$  PERTURBATION SERIES IN  $\alpha$

e.g.  $e^-$  MAGNETIC MOMENT

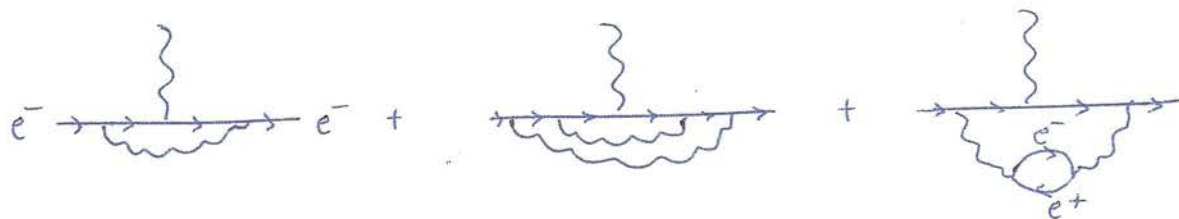
$\mu = g \cdot \underbrace{\frac{1}{2}}_{\text{SPIN}} \cdot \left( \frac{e}{2m_e} \right)$  BOHR MAGNETON

• FOR DIRAC PARTICLE  $g = 2$



- DUE TO VIRTUAL  $\gamma$  PROCESSES (VACUUM POLARIZATION, ...) INT 6  
 $g \neq 2$  ↓  
QUANTUM FIELD THEORY

$\frac{g-2}{2}$  :  $e^-$  ANOMALOUS MAGNETIC MOMENT



PERTURBATION SERIES

$$\frac{g-2}{2} = 0.5 \left( \frac{\alpha}{\pi} \right) - 0.32848 \left( \frac{\alpha}{\pi} \right)^2 + O\left( \frac{\alpha}{\pi} \right)^3$$

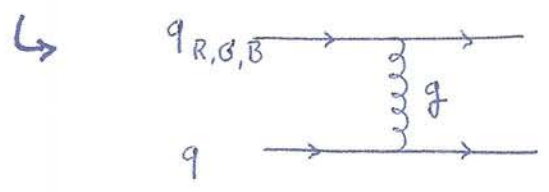
↑  
0.00116

$$\left( \frac{g-2}{2} \right)^{\text{THEORY}} = (11\,659\,184.0 \pm 5.9) \times 10^{-10}$$

$$\left( \frac{g-2}{2} \right)^{\text{EXP}} = (11\,659\,208.9 \pm 6.3) \times 10^{-10}$$

THEORY  $\leftrightarrow$  EXP  
 (2012) 30 DEVIATION

\* STRONG INTERACTIONS



8 gluons (  $R\bar{B}$ ,  $R\bar{G}$ ,  $B\bar{R}$ ,  $B\bar{G}$ ,  $G\bar{R}$ ,  $G\bar{B}$ ,  $R\bar{R} - G\bar{G}$ ,  $R\bar{R} + G\bar{G} - 2B\bar{B}$  )

SU(3) GAUGE SYMMETRY : THEORY IS INVARIANT WHEN QUARK CHANGES ITS COLOR

MASSLESS → PROPAGATOR  $\sim \frac{1}{q^2}$

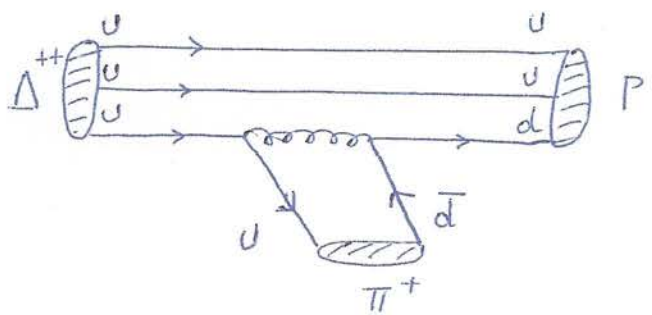
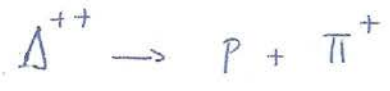
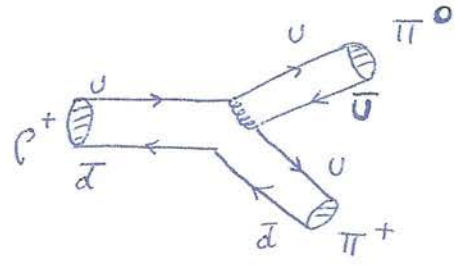


$\frac{1}{r}$  POTENTIAL AT SHORT DISTANCES

$$V(r) = \frac{a}{r} + \underbrace{br}_{\text{CONFINEMENT POTENTIAL}}$$

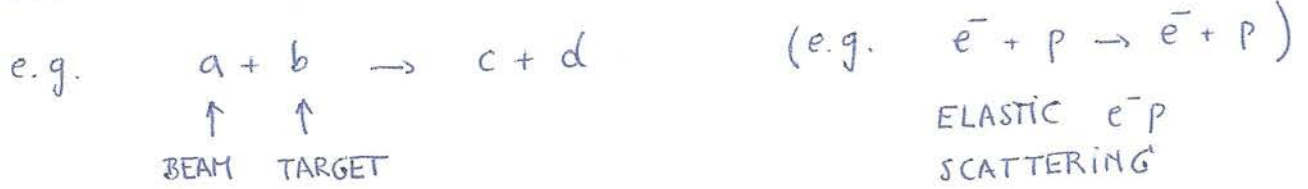
↑  
1g EXCHANGE POTENTIAL AT LARGE DISTANCES

↳ STRONG DECAYS



\* INTERACTION CROSS SECTION

STRENGTH OF SCATTERING PROCESS

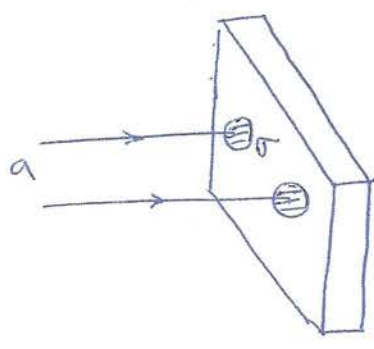


$\sigma$  : REACTION CROSS SECTION (FOR EACH TARGET PARTICLE)  
 $\equiv \frac{\text{TRANSITION PROBABILITY PER UNIT TIME : } W}{\text{INCIDENT FLUX : } \phi}$  (DIMENSION SURFACE).

$\rightarrow$  INCIDENT FLUX  $\phi = n_a v_i$   
 $\uparrow \quad \uparrow$   
 DENSITY OF PARTICLES IN BEAM RELATIVE VELOCITY OF BEAM (w.r.t TARGET)

$\rightarrow$  # REACTIONS PER UNIT TIME & PER UNIT SURFACE

$N = \phi \cdot (n_b \sigma dx)$



$dx$  (TARGET THICKNESS)

$n_b$  : DENSITY OF TARGET

$n_b dx$  : # TARGET PARTICLES / UNIT SURFACE



→  $W$ : REACTION RATE (PER TARGET PARTICLE)

- NON-RELATIV. (FERMI'S GOLDEN RULE)

$$W = \frac{2\pi}{\hbar} |M_{fi}|^2 \cdot \rho_f$$

↑  
MATRIX ELEMENT  
 $\langle f | \hat{O} | i \rangle$

←  
PHASE SPACE  
OF FINAL STATE  
(ENERGY DENSITY)

- RELATIVISTICALLY

CALCULATE  $M_{fi}$  FROM FEYNMAN DIAGRAM

$$\rho_f = \frac{d^3 \bar{p}_f}{(2\pi)^3 2E_f} \quad \text{FOR 1 PARTICLE IN FINAL STATE}$$

$$\rho_f = \frac{d^3 \bar{p}_c}{(2\pi)^3 2E_c} \frac{d^3 \bar{p}_d}{(2\pi)^3 2E_d} \quad \text{FOR 2 PARTICLE FINAL STATE}$$

→ ENERGY - MOMENTUM IS CONSERVED

$$\text{FACTOR } (2\pi)^4 \delta^4 (p_a + p_b - p_c - p_d)$$

→ SPIN STATES

$a$ : SPIN  $s_a \Rightarrow (2s_a + 1)$  STATES

AVERAGE OVER INITIAL SPINS

& SUM OVER FINAL SPINS

$$\frac{1}{(2s_a + 1)(2s_b + 1)} \sum_{s_a, s_b, s_c, s_d}$$

→ UNITS OF  $\sigma$

barns	$1 \text{ barn} = 10^{-28} \text{ m}^2$
millibarn	$1 \text{ mb} = 10^{-3} \text{ b}$
microbarn	$1 \mu\text{b} = 10^{-6} \text{ b}$
nanobarn	$1 \text{ nb} = 10^{-9} \text{ b}$
picobarn	$1 \text{ pb} = 10^{-12} \text{ b}$

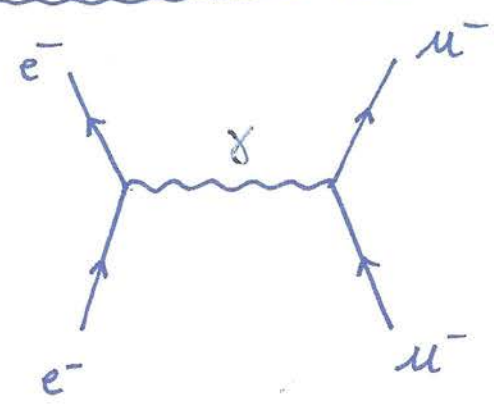
EXPRESS  $\sigma$  IN  $\text{GeV}^{-2}$

& CONVERT BY  $(\hbar c)^2 = (0.197 \text{ GeV}^2 \text{ fm}^2)$

+ express  $\text{fm}^2$  IN b

$$1 \text{ fm}^2 = 10^{-2} \text{ b}$$

\* EXAMPLE : " SPINLESS " ELECTRON - MUON SCATTERING



⇒ e<sup>-</sup> & μ<sup>-</sup> ARE SPIN 1/2 DIRAC PARTICLES

LET'S FIRST WORK IT OUT FOR THE CASE OF SPIN 0 PARTICLES (e.g. π<sup>-</sup>, π<sup>+</sup>)

⇒ SPIN - 0 PARTICLE SATISFIES KLEIN - GORDON EQ.

$$(\partial_\mu \partial^\mu + m^2) \phi = 0 \quad \left( \partial_\mu \equiv \frac{\partial}{\partial x^\mu} \right)$$

THIS FIELD EQ. FOLLOWS FROM LAGRANGIAN

$$L = \int d^3\vec{x} \mathcal{L}, \quad \mathcal{L}_{KG} = (\partial_\mu \phi)^\dagger (\partial^\mu \phi) - m^2 \phi^\dagger \phi$$

⇓  
USE EULER - LAGRANGE EQUATIONS

$$\frac{\partial \mathcal{L}}{\partial \phi} - \partial_\mu \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} = 0 \quad \Rightarrow (\partial_\mu \partial^\mu + m^2) \phi = 0$$

& LIKEWISE FOR  $\phi^\dagger$



INTERACTION OF CHARGED PARTICLE (CHARGE  $e$ )  
WITH ELECTROMAGNETIC FIELD

$$p^\mu \rightarrow p^\mu - e A^\mu$$

FOR  $\phi(x) = \phi_0 e^{-i p \cdot x}$  PLANE WAVE

$$\partial^\mu \rightarrow \partial^\mu + i e A^\mu$$

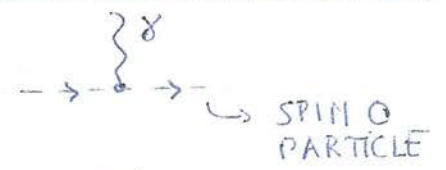
 MINIMAL SUBSTITUTION

↓  
PERFORM THIS SUBSTITUTION IN  $\mathcal{L}_{KG}$

$$\begin{aligned} \mathcal{L}_{KG} + \mathcal{L}_{INT} &= \left[ (\partial_\mu + i e A_\mu) \phi \right]^\dagger (\partial^\mu + i e A^\mu) \phi - m^2 \phi^\dagger \phi \\ &= (\partial_\mu \phi^\dagger) (\partial^\mu \phi) - m^2 \phi^\dagger \phi \\ &\quad - i e A_\mu \phi^\dagger \partial^\mu \phi + i e (\partial_\mu \phi)^\dagger \phi A^\mu \\ &\quad + e^2 \phi^\dagger \phi A_\mu A^\mu \end{aligned}$$

$$\mathcal{L}_{INT} = - e \left[ i \phi^\dagger (\partial^\mu \phi) - i (\partial^\mu \phi)^\dagger \phi \right] A_\mu + e^2 \phi^\dagger \phi A_\mu A^\mu$$

TERM LINEAR IN  $A^\mu$



TERM QUADRATIC IN  $A^\mu$

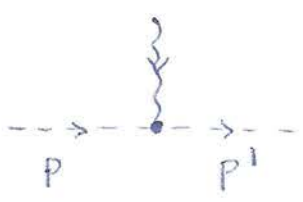


- TERM LINEAR IN  $A^\mu$

$$\mathcal{L}_{\text{INT}} = - e J^\mu A_\mu$$

$J^\mu$ : ELECTROMAGNETIC CURRENT

$$J^\mu(x) = i \phi^\dagger (\partial^\mu \phi) - i (\partial^\mu \phi^\dagger) \phi$$

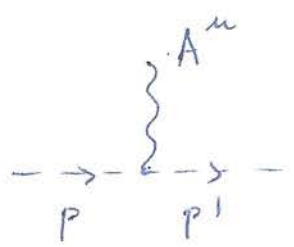


INITIAL  $\phi_i \sim e^{-i p \cdot x}$   
 FINAL  $\phi_f^\dagger \sim e^{+i p' \cdot x}$

IN MOMENTUM SPACE

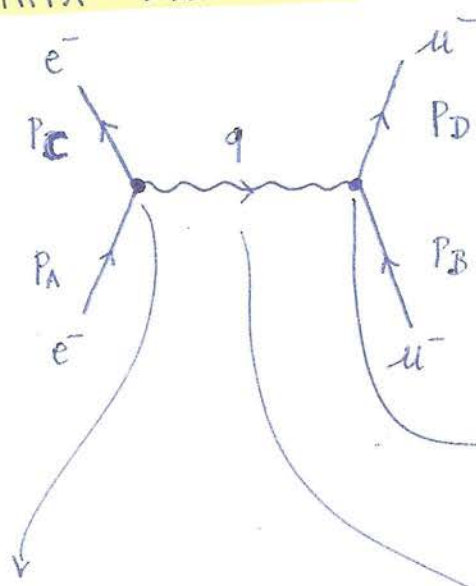
$$\langle p' | J^\mu(0) | p \rangle = (p + p')^\mu$$

- FEYNMAN RULE (LOWEST ORDER:  $i\mathcal{L}$ )



$$- i e (p + p')^\mu$$

⇒ **MATRIX ELEMENT** FOR SPINLESS  $e^-u^-$  SCATTERING



$e^-$  HAS CHARGE  $(-e)$   
 $u^-$

$$[-i(-e)(p_A + p_C)^\mu]$$

$$[-i(-e)(p_B + p_D)^\nu]$$

PHOTON PROPAGATOR  
 (LORENZ GAUGE)

$$-\frac{i g_{\mu\nu}}{q^2}$$

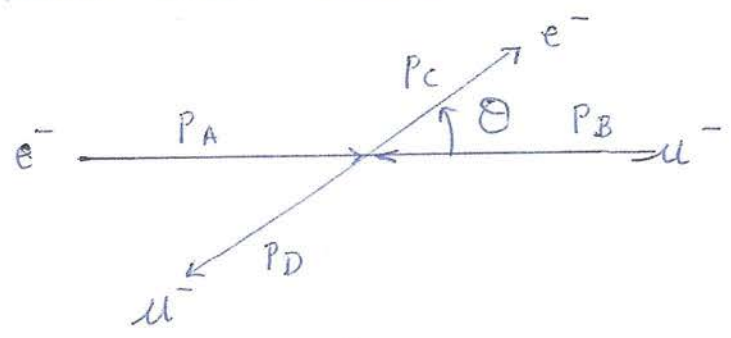
$$\mathcal{M}_{AB \rightarrow CD} = [ +ie(p_A + p_C)^\mu ] \left( -\frac{i g_{\mu\nu}}{q^2} \right) [ ie(p_B + p_D)^\nu ]$$

↑  
MATRIX ELEMENT

$$\mathcal{M}_{AB \rightarrow CD} = \frac{ie^2}{q^2} (p_A + p_C)^\mu (p_B + p_D)_\mu$$

⇒ CROSS SECTION

- IN TOTAL C.M. (CENTER-OF-MASS) FRAME



$$P_A = (E_A, \vec{P}_A)$$

$$P_B = (E_B, \vec{P}_B) \quad \text{IN C.M. } \vec{P}_B = -\vec{P}_A, E_B = E_A$$

$$P_C = (E_C, \vec{P}_C)$$

TOTAL ENERGY IS CONSERVED

$$P_D = (E_C, -\vec{P}_D)$$

$$E_C = E_A$$

TOTAL ENERGY INITIALLY  $E_A + E_B = 2E_A$

INVARIANT  $(P_A + P_B)^2 \equiv s = (2E_A)^2 - (\vec{P}_A + \vec{P}_B)^2$

$$\boxed{E_A = \frac{\sqrt{s}}{2}}$$

$$E_C = \sqrt{s}/2$$

s : MANDELSTAM INVARIANT

LET'S CONSIDER HIGH BEAM ENERGIES

↳  $|\vec{P}_A| \gg m_e, m_\mu \Rightarrow$  WE CAN NEGLECT MASSES

$$P_A \left( \frac{\sqrt{s}}{2}, 0, 0, \frac{\sqrt{s}}{2} \right)$$

$$E_A \approx |\vec{P}_A|$$

$$P_B \left( \frac{\sqrt{s}}{2}, 0, 0, -\frac{\sqrt{s}}{2} \right)$$

$$P_C \left( \frac{\sqrt{s}}{2}, \frac{\sqrt{s}}{2} \sin \theta, 0, \frac{\sqrt{s}}{2} \cos \theta \right)$$

$$P_D \left( \frac{\sqrt{s}}{2}, -\frac{\sqrt{s}}{2} \sin \theta, 0, -\frac{\sqrt{s}}{2} \cos \theta \right)$$

$$\mathcal{K}_{AB \rightarrow CD} = \frac{ie^2}{q^2} (P_A + P_C)^{\mu} (P_B + P_D)_{\mu}$$

$$\Rightarrow (P_A + P_C)^{\mu} (P_B + P_D)_{\mu}$$

$$= (\sqrt{s}) \cdot (\sqrt{s})$$

$$- \left( \frac{\sqrt{s}}{2} \sin \theta \right) \cdot \left( -\frac{\sqrt{s}}{2} \sin \theta \right)$$

$$- \left( \frac{\sqrt{s}}{2} (1 + \cos \theta) \right) \cdot \left( -\frac{\sqrt{s}}{2} (1 + \cos \theta) \right)$$

$$= s \left( 1 + \frac{1}{4} \sin^2 \theta + \frac{1}{4} (1 + \cos \theta)^2 \right)$$

$$= \frac{s}{4} (6 + 2 \cos \theta)$$

$$= \frac{s}{2} (3 + \cos \theta)$$

$$\Rightarrow q^2 = (P_A - P_C)^2 = -\frac{s}{4} (1 - \cos \theta)^2 - \frac{s}{4} \sin^2 \theta$$

$$= -\frac{s}{2} \sin^2 \theta$$

$$\left\| \mathcal{K}_{AB \rightarrow CD} = -ie^2 \frac{(3 + \cos \theta)}{(1 - \cos \theta)} = -\frac{ie^2}{2} \frac{(3 + \cos \theta)}{\sin^2 \frac{\theta}{2}} \right.$$



$$d\sigma = \frac{1}{\Phi} \cdot W$$

$\uparrow$  INITIAL FLUX       $\uparrow$  TRANSITION PROB / UNIT TIME

$$\Phi = \frac{|\bar{P}_A|}{E_A} + \frac{|\bar{P}_B|}{E_B}$$

TOTAL CURRENT (INITIAL) DENSITY

$$= 2$$

$$W = \frac{1}{(2E_A)(2E_B)} \frac{d^3 \bar{P}_C}{(2\pi)^3 2E_C} \cdot \frac{d^3 \bar{P}_D}{(2\pi)^3 2E_D} \cdot (2\pi)^4 \delta^4(P_A + P_B - P_C - P_D)$$

$\cdot |\mathcal{M}_{AB \rightarrow CD}|^2$

$$d\sigma = \frac{1}{2s} \cdot \frac{1}{(2\pi)^2} \frac{d^3 \bar{P}_C}{4E_C E_D} \delta \left( \underbrace{P_A^0 + P_B^0}_{\sqrt{s}} - \sqrt{\bar{P}_C^2 + m_e^2} - \sqrt{\bar{P}_D^2 + m_e^2} \right)$$

$\uparrow$   $|\bar{P}_D| = |\bar{P}_C|$

$\cdot |\mathcal{M}_{AB \rightarrow CD}|^2$

$$= \frac{1}{2s^2 (2\pi)^2} \cdot d\Omega \cdot d|\bar{P}_C| |\bar{P}_C|^2 \underbrace{\delta(\sqrt{s} - 2|\bar{P}_C|)}_{\frac{1}{2} \delta(|\bar{P}_C| - \frac{\sqrt{s}}{2})} |\mathcal{M}_{AB \rightarrow CD}|^2$$

$$\left( \frac{d\sigma}{d\Omega} \right)^{c.m.} = \frac{1}{2s^2 (2\pi)^2} \cdot \left( \frac{\sqrt{s}}{2} \right)^2 \cdot \frac{1}{2} \cdot |\mathcal{M}_{AB \rightarrow CD}|^2$$

$$\left(\frac{d\sigma}{d\Omega}\right)^{\text{c.H.}} = \frac{1}{4s (4\pi)^2} \cdot e^4 \frac{(3 + \cos\theta)^2}{(1 - \cos\theta)^2}$$

↓ INTRODUCE  $\alpha \equiv \frac{e^2}{4\pi}$  FINE STRUCTURE CONST.

$$\left(\frac{d\sigma}{d\Omega}\right)^{\text{c.H.}} = \frac{\alpha^2}{4s} \cdot \frac{(3 + \cos\theta)^2}{(1 - \cos\theta)^2} = \frac{\alpha^2}{16s} \frac{(3 + \cos\theta)^2}{\sin^4\theta/2}$$

↑  
STRONG FORWARD PEAKING

$\sim \frac{1}{\sin^4\theta/2}$  RUTHERFORD FORMULA

NUMERICAL EXAMPLE:  $\theta^{\text{c.H.}} = 90^\circ$ ,  $\sqrt{s} = 100 \text{ GeV}$

$$\left.\frac{d\sigma}{d\Omega}\right|_{\theta=90^\circ} = \frac{g\alpha^2}{4s} = \frac{g}{4} \left(\frac{1}{137}\right)^2 \cdot \frac{1}{10^4} \text{ GeV}^{-2}$$

$$= \frac{g}{4} \left(\frac{1}{137}\right)^2 \cdot \frac{(0.197)^2}{10^4} \text{ fm}^2$$

$\underbrace{\hspace{2cm}}_{10^{-2} \text{ b}}$

$$= \frac{g}{4} \left(\frac{197}{137}\right)^2 \cdot \underbrace{10^{-12}}_{\text{pb}} \text{ b}$$

$\approx 4.65$