# Practice Exam <br> Theoretical Physics 6a (QFT): SS 2019 

08.07.2019

## Exercise 1. $e^{+} e^{-} \rightarrow \phi^{+} \phi^{-}$(50 points)

In the lecture, we derived the differential cross section $(d \sigma / d \Omega)$ for the $e^{+} e^{-} \rightarrow \mu^{+} \mu^{-}$in the collider frame, assuming the high energy limit $\left(\sqrt{s} \gg m_{\mu}\right)$.
For this exercise, consider the case when the muons (spin $1 / 2$ ) are replaced by spinless bosons $\phi$ (spin $0)$.
(a)(20 points)Determine the possible Feynman diagrams at leading order and calculate the corresponding S-matrix element ( $S_{f i}$ ) for the $e^{+} e^{-} \rightarrow \phi^{+} \phi^{-}$process using Feynman rules.
Hint: Notice that the scalar QED interaction needs to be considered in this case.
(b)(25 points) Calculate the unpolarized differential cross-section $d \sigma / d \Omega$ for this process in the the center of mass frame (collider frame).
(c)(05 points) Using the result obtained in the previous item, determine the total cross section.

## Exercise 2. Conformal Symmetry (15 points)

Consider the Lagrangian for scalar $\phi^{4}$ theory:

$$
\begin{equation*}
\mathcal{L}=\frac{1}{2} \partial_{\mu} \phi \partial^{\mu} \phi-\frac{1}{2} m^{2} \phi^{2}-\frac{g}{4!} \phi^{4} \tag{1}
\end{equation*}
$$

a) (5 p.) Calculate the energy momentum tensor

$$
\begin{equation*}
T^{\mu \nu}=\frac{\partial \mathcal{L}}{\partial\left(\partial_{\mu} \phi\right)} \partial^{\nu} \phi-g^{\mu \nu} \mathcal{L}, \tag{2}
\end{equation*}
$$

and show explicitly by using the equation of motion for the field $\phi$ that

$$
\begin{equation*}
\partial_{\mu} T^{\mu \nu}=0 . \tag{3}
\end{equation*}
$$

From which symmetry transformation follows the conservation law of the energy momentum tensor?
b) (5 p.) Calculate the trace $T_{\mu}^{\mu}$ of the energy-momentum tensor and show that it can be written as a divergence in the case of $m=0$.
c) (5 p.) Consider a infinitesimal scaling transformation

$$
\begin{equation*}
x^{\prime}=(1-\epsilon) x . \tag{4}
\end{equation*}
$$

This transformation induces a transformation law for the scalar field as:

$$
\begin{equation*}
\phi^{\prime}\left(x^{\prime}\right)=\left(1+\epsilon d_{\phi}\right) \phi(x), \tag{5}
\end{equation*}
$$

where $d_{\phi}$ is the scaling parameter of the field $\phi$. Show that the variation of the field $\phi \mathrm{s}$ given by:

$$
\begin{equation*}
\delta \phi=\phi^{\prime}(x)-\phi(x)=\epsilon\left(d_{\phi}+x^{\mu} \partial_{\mu}\right) \phi \tag{6}
\end{equation*}
$$

## Exercise 3. 1-loop correction to the propagator in Yukawa theory +10 Bonus points)

Consider the interaction between a scalar field $\phi$ (with mass $M$ ) and a spin $1 / 2$ field $\psi$ (with mass $m$ ) described by the Lagrangian:

$$
\begin{equation*}
\mathcal{L}=\frac{1}{2}\left(\partial_{\mu} \phi\right)\left(\partial^{\mu} \phi\right)-\frac{1}{2} M^{2} \phi^{2}+\bar{\psi}\left(i \gamma^{\mu} \partial_{\mu}-m\right) \psi-\lambda \bar{\psi} \psi \phi \tag{7}
\end{equation*}
$$

where $\lambda$ is a coupling constant.
a) ( 5 p.) Derive the Feynman rule corresponding to the interaction term.
b) (10 p.) The 1-loop correction to the scalar propagator induced by a fermion loop is presented in Fig. 3. Use the Feynman rules to derive the invariant amplitude $\mathcal{M}$ for this diagram, using the momentum labels as indicated on the figure.


Figure 1: One-loop correction to the scalar propagator in Yukawa theory, given by the Lagrangian in Eq. (7).
c) (10 p.) Use the Feynman parameterization, and perform the one-loop integral using dimensional regularization (and using the formulas given at the end). Show that the result can be expressed as:

$$
\begin{equation*}
\mathcal{M}^{1-\text { loop }}=i \frac{4(d-1) \lambda^{2} \mu^{4-d}}{(4 \pi)^{d / 2}} \Gamma(1-d / 2) \int_{0}^{1} d x \frac{1}{\left[m^{2}-p^{2} x(1-x)\right]^{1-d / 2}} \tag{8}
\end{equation*}
$$

where $d$ denotes the dimensionality of space-time, and $\mu$ is some arbitrary scale to keep the coupling $\lambda$ dimensionless.
d) (10 p.) The scalar counterterms $(C T)$ that have to be added to the diagram of Fig. 2 correspond with the Feynman rule:

$$
\begin{equation*}
\mathcal{M}^{C T}=i\left[p^{2} \delta_{\phi}-M^{2}\left(\delta_{M}+\delta_{\phi}\right)\right] \tag{9}
\end{equation*}
$$

where $\delta_{\phi}$ is the counterterm for the field $\phi$, and $\delta_{M}$ is the counterterm for the scalar squared mass $M^{2}$.
Defining $\varepsilon \equiv 2-d / 2$, expand the above result for the invariant amplitude $\mathcal{M}$ in $\varepsilon$ to extract the pole term in $1 / \varepsilon$. Use the $M S$ subtraction scheme, i.e. absorb only the divergent parts, and determine the $M S$ expressions for the counterterms $\delta_{\phi}$ and $\delta_{M}$.
e) (10 Bonus p.) The renormalized propagator of the scalar field is given by

$$
\begin{equation*}
\frac{i}{p^{2}-M^{2}-\Sigma_{R}\left(p^{2}\right)} \tag{10}
\end{equation*}
$$

with the renormalized self-energy $\Sigma_{R}\left(p^{2}\right)=i\left(\mathcal{M}^{1-\text { loop }}+\mathcal{M}^{C T}\right)$. Using the above result for the invariant amplitude, what is the expression for $\Sigma_{R}\left(p^{2}\right)$ in the $\overline{M S}$ scheme? You do not need to perform the Feynman parameter integral. What is the expression for the difference between the pole value $\left(M_{P}^{2}\right)$ and the $\overline{M S}$ value $\left(M_{\overline{M S}}^{2}\right)$ of the squared scalar mass ?
Note: This difference determines the shift in the Higgs mass ( $M$ ) due to the heavy (mass m) top-quark loop.

## Useful Formulas

The formula for the cross section for the process $a+b \rightarrow c+d$ is given by:

$$
\begin{equation*}
d \sigma=\frac{1}{\left(2 E_{a}\right)\left(2 E_{b}\right) v_{r e l}} \frac{d^{3} \vec{p}_{c}}{(2 \pi)^{3}\left(2 E_{c}\right)} \frac{d^{3} \vec{p}_{d}}{(2 \pi)^{3}\left(2 E_{d}\right)}(2 \pi)^{4} \delta^{4}\left(p_{a}+p_{b}-p_{c}-p_{d}\right)|\mathcal{M}|^{2} \tag{11}
\end{equation*}
$$

where the four-momenta for each particle are given by $p_{i}=\left(E_{i}, \vec{p}_{i}\right)$, where $v_{r e l}$ stands for the incident flux, and $\mathcal{M}$ is the invariant amplitude.

$$
\begin{gather*}
\frac{1}{D_{1} D_{2}}=\int_{0}^{1} d x \frac{1}{\left[(1-x) D_{1}+x D_{2}\right]^{2}}  \tag{12}\\
\int \frac{d^{d} q}{(2 \pi)^{d}} \frac{1}{\left(q^{2}-\Delta+i \epsilon\right)^{n}}=\frac{i}{(4 \pi)^{d / 2}} \frac{\Gamma(n-d / 2)}{\Gamma(n)} \frac{(-1)^{n}}{\Delta^{n-d / 2}}  \tag{13}\\
\int \frac{d^{d} q}{(2 \pi)^{d}} \frac{q^{\mu} q^{\nu}}{\left(q^{2}-\Delta+i \epsilon\right)^{n}}=g^{\mu \nu} \frac{i}{(4 \pi)^{d / 2}} \frac{\Gamma(n-1-d / 2)}{2 \Gamma(n)} \frac{(-1)^{n-1}}{\Delta^{n-1-d / 2}}  \tag{14}\\
\Gamma(z+1)=z \Gamma(z)  \tag{15}\\
\Gamma(-1+\varepsilon)=-1\left\{\frac{1}{\varepsilon}+1-\gamma_{E}+\mathcal{O}(\varepsilon)\right\} \tag{16}
\end{gather*}
$$

with Euler constant $\gamma_{E} \approx 0.577$.

