

## Cosmology and General Relativity: HW 5

Turn in to "Wittig" (Hausaufgaben) mailbox in KPH by **noon, 27 June 2019**

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**Problem 1 Maximally symmetric spaces:** The easiest example of a maximally symmetric space is  $\mathbb{R}^n$  with the flat Euclidean metric. Its isometries are the rotation and translations in  $n$  dimensions. In the neighborhood of some fixed point  $p$  the translations move this point along an axis, and since there are  $n$  independent axes, there are  $n$  possible translations. Rotations leave  $p$  invariant, but transform one axis through  $p$  and into another one. Since there are  $n$  axes, we can move each into  $n - 1$  other axes. After considering double-counting, there are  $\frac{1}{2}n(n - 1)$  independent rotations. In total, there are  $\frac{1}{2}n(n + 1)$  independent symmetries in  $\mathbb{R}^n$ . Because we considered the local effect of the symmetries these arguments should also apply for curved spaces. If the manifold is pseudo-Riemannian, some rotations will be boosts, but the counting is the same. Since the number of isometries is the number of linearly independent Killing Vector fields, we can call an  $n$ -dimensional manifold with  $\frac{1}{2}n(n + 1)$  Killing vectors *maximally symmetric*. For a maximally symmetric manifold, the curvature is the same everywhere and in every direction. This we can classify maximally symmetric manifolds by their Ricci curvature scalar  $R$ , their dimension  $n$ , their metric signature, and perhaps some additional information on the global topology (e.g., the size of the manifold). For a given dimension and signature, we can classify maximally symmetric spaces (ignoring global topology) into the basic categories of  $R$  being positive, negative, or zero; since the magnitude of  $R$  corresponds to an overall scaling of the space. For Riemannian manifolds, the flat (i.e.,  $R = 0$ ) maximally symmetric spaces correspond to planes (or their high dimensional generalizations).

We want to consider the two-dimensional hyperboloid  $H^2$ . There are several ways to represent  $H^2$  (which has the same topology as  $\mathbb{R}^2$ ). We choose the *Poincaré half-plane*, being defined on the upper half-plane of  $\mathbb{R}^2$  with the metric

$$ds^2 = \frac{a^2}{y^2} (dx^2 + dy^2) .$$

- Compute the lengths of vertical line segments ( $x = \text{const}$ ) from  $y_1$  to  $y_2$ . What happens for  $y \rightarrow 0$ ? 1pt
- Calculate the Christoffel symbols and show that the geodesics satisfy

$$(x - x_0)^2 + y^2 = \ell^2 ,$$

for some constants  $x_0$  and  $\ell$ . What do these geodesics look like? 2pt

- Calculate the Riemannian tensor and Ricci tensor. Show that the Ricci curvature scalar is constant on the manifold. 2.5pt

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**Problem 2 Killing vector fields on the 2-sphere:** Consider the 2-sphere  $S^2$  with the usual induced metric,  $ds^2 = d\theta^2 + \sin^2 \theta d\varphi^2$ .

- a. Write down the three independent Killing equations

$$\nabla_\mu X_\nu + \nabla_\nu X_\mu = 0 ,$$

for a vector field  $X = X^\theta \partial_\theta + X^\varphi \partial_\varphi$ . 1pt

- b. Show that  $X_\theta$  is independent of  $\theta$ , i.e. that you can also write  $X_\theta(\theta, \varphi) = f(\varphi)$ . Perform this substitution in the Killing equations and show that:

$$X_\varphi = -F(\varphi) \sin \theta \cos \theta + g(\theta) ,$$

where  $F(\varphi)$  is the primitive of  $f(\varphi)$ , so  $F'(\varphi) = f(\varphi)$ , and  $g(\varphi)$  is an integration constant. 1.5pt

- c. Insert the result of the previous item in the last Killing equation and — using separation of variables — show that one obtains:

$$\frac{dg}{d\theta} - 2g(\theta) \cot \theta = C , \tag{1}$$

$$\frac{df}{d\varphi} + F(\varphi) = -C , \tag{2}$$

where  $C$  is a constant. 1.5pt

- d. Determine  $g(\theta)$  by integrating equation (1). Differentiate equation (2) and show that  $f$  is harmonic. Give the general solution.

*Hint: The result has to be of the form:* 1pt

$$g(\theta) = (C_1 - C \cot \theta) \sin^2 \theta, \quad f(\varphi) = A \sin \varphi + B \cos \varphi.$$

- e. Combine the results and show that a general Killing vector on  $S^2$  has the following form: 1pt

$$X = A(\sin \varphi \partial_\theta + \cos \varphi \cot \theta \partial_\varphi) + B(\cos \varphi \partial_\theta - \sin \varphi \cot \theta \partial_\varphi) + C_1 \partial_\varphi . \tag{3}$$

- f. Identify the three basis vectors of the Killing vector in equation (3) with the angular momentum  $L_i = \sum_{j,k} \epsilon_{ijk} x_j \partial_k$ . Is  $S^2$  a maximally symmetric space? 2pt

**Problem 3 Spectral Red Shift:** The spherically symmetric solution to Einstein's equations for a massive object of mass  $m$  (here  $1 = c$ ) is given by:

$$ds^2 = -A(r)dt^2 + B(r)dr^2 + r^2(d\vartheta^2 + \sin^2\vartheta d\varphi^2)$$

with

$$A(r) = \left(1 - \frac{2GM}{r}\right) = \frac{1}{B(r)} .$$

Assume that an *emitter* sends a signal at a fixed point  $(r_E, \vartheta_E, \varphi_E)$  and that the signal propagates along the null geodesic and is detected by the *receiver* at a fixed point  $(r_R, \vartheta_R, \varphi_R)$ . If  $t_E$  is the time coordinate of emission and  $t_R$  the time coordinate of reception, the signal propagates from event  $(t_E, r_E, \vartheta_E, \varphi_E)$  to  $(t_R, r_R, \vartheta_R, \varphi_R)$ .

- a. Draw the spacetime diagram and illustrate both events. 1pt
- b. Let  $\lambda$  be the affine parameterization of the null geodesic with  $\lambda_{E/R}$  at the points of emission/reception. Show that the following holds ( $i, j = 1, 2, 3$ ):

$$\frac{dt}{d\lambda} = \left[ \left(1 - \frac{2GM}{r}\right)^{-1} g_{ij} \frac{dx^i}{d\lambda} \frac{dx^j}{d\lambda} \right]^{1/2} .$$

2pt

- c. Take the above even to argue that

$$\Delta t_E = \Delta t_R \quad \text{with} \quad \Delta t = t^{(2)} - t^{(1)}$$

corresponds to the difference of the coordinate times of signals 1 and 2. 1pt

- d. A clock at the position of an observer measures the proper time  $\tau$  instead of the coordinate time  $t$ . Find the relation among the two time parameterization and conclude:

$$\frac{\Delta\tau_R}{\Delta\tau_E} = \left[ \frac{1 - 2MG/r_R}{1 - 2MG/r_e} \right]^{1/2} .$$

1pt

- e. Now assume that the emitter is pulsed with frequency  $\nu_E = \frac{n}{\Delta\tau_E}$ , i.e.  $n$  pulses per proper time interval  $\Delta\tau_E$  are emitted. Analogous relations hold for the receiver. Find the relation among the two frequencies  $\nu_E/\nu_R$ . Expand the relation for  $r_E, r_R \gg 2GM$  and discuss what happens if the emitter (receiver) is closer to the massive object than the receiver (emitter). 1.5pt