

# Problem Sheet 2

for the course  
„Introduction to Lattice Gauge Theory“  
Summer 2019

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## 1. Reading List

1. Smit, chapter 3 (p. 32–55)
2. Parisi, chapters 3 and 4 (p. 22–66)

## 2. The Hopping-Parameter Expansion

- (a) Consider the theory of a free scalar field  $\phi$  with mass  $m_0$  and split the action into a local term and a term coupling neighboring lattice points. Perform a rescaling  $\phi = \sqrt{2\kappa}\varphi$  and determine the value of  $\kappa$  as a function of  $m_0$  for which the coefficient of the local term becomes equal to one.
- (b) In the following we assume that  $\kappa$  has the value determined in the preceding problem. Which critical value  $\kappa_c$  of  $\kappa$  corresponds to the massless case?
- (c) Consider now a correlation function of the form  $\langle \varphi_x \varphi_y \rangle$  and expand the corresponding path integral into a power series in  $\kappa$ . Which terms give a non-vanishing contribution? Interpret the result graphically.

## 3. From the $\phi^4$ -Theory to the Ising Model

- (a) We now add an interaction  $V_I(\phi) = \frac{\lambda_0}{4!} \sum_x \phi_x^4$  to the theory from the preceding problem. For which values of  $\kappa$  and  $\lambda$  as a function of  $m_0$  and  $\lambda_0$  does the local term take the form  $\varphi_x^2 + \lambda(\varphi_x^2 - 1)^2$  ?
- (b) Consider now the limit  $\lambda \rightarrow \infty$  and show that in this limit the theory can be written as an Ising model with partition function

$$Z = \sum_{\{\sigma=\pm 1\}} e^{-\beta H(\sigma)}$$

and Hamiltonian

$$H = -J \sum_{(i,j)} \sigma_i \sigma_j$$

where the sum over  $(i, j)$  run over all pairs of neighboring lattice sites.

4. *Solution of the Ising model in one dimension and in infinitely many dimensions*

- (a) Write the partition function for the one-dimensional Ising model using the transfer matrix (here literally a  $2 \times 2$  matrix). By diagonalizing the transfer matrix, derive the exact solution for the free energy per site  $\frac{F}{L} = -\frac{1}{\beta L} \log Z$  in the thermodynamic limit  $L \rightarrow \infty$ .
- (b) In the opposite case of infinitely many dimensions,  $D \rightarrow \infty$ , explain why we can replace the sum of the neighboring spins by the expectation value of the spin:  $\sum_{(i,j)} \sigma_j \rightarrow 2D \langle \sigma_i \rangle$ . Use this ansatz to derive an equation which the magnetization  $m = \langle \sigma_i \rangle$  has to satisfy.
- (c) If we use the infinite-dimensional ansatz also in a finite number  $D$  of dimensions, we obtain an approximation (the mean-field approximation). Under which conditions does the resulting equation have how many solutions? Interpret the result graphically and physically.

5. *Low-Temperature Expansion of the Ising Model*

- (a) Find the ground state of the Ising model in a volume  $V = L^D$  with periodic boundary conditions, and determine its energy.
- (b) Considering now states that differ from the ground state by flipping one and two spins, respectively, determine how many such states there are in each case, and determine their energies.
- (c) Use the results of the preceding two questions to expand the partition function  $Z$  into a power series in a suitable variable  $\lambda$  (which is to satisfy  $\lambda \rightarrow 0$  for  $\beta \rightarrow \infty$ ).
- (d) Explain why the free energy per site,  $\frac{F}{V}$ , is independent of  $V$ .

6. *High-Temperature Expansion of the Ising Model*

- (a) Show the identity

$$e^{\beta s s'} = \cosh \beta (1 + s s' \tanh \beta)$$

for  $s, s' = \pm 1$ .

- (b) Use this identity to rewrite the partition function of the Ising model as

$$Z = (\cosh(\beta J))^V \sum_{\{n_{ij}=0,1\}} \kappa^{\sum_{(i,j)} n_{ij}} \sum_{\{\sigma\}} \prod_{(i,j)} (\sigma_i \sigma_j)^{n_{ij}}$$

with  $\kappa = \tanh(\beta J)$ .

- (c) Consider the contribution of an individual spin  $\sigma_i$  to the partition function and show that it contributes a factor 2 to  $\sum_{\{\sigma\}} \prod_{(i,j)} (\sigma_i \sigma_j)^{n_{ij}}$  if  $\sum_k n_{lk}$  is even, whereas the partition function vanishes otherwise.
- (d) Use the result of the preceding question to rewrite  $Z$  as a sum over configurations of only the auxiliary variables  $n_{ij}$ , and interpret the result geometrically.