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A simpler and more elegant scheme can be constructed if we allow non-integral values for the charges. We can dispense entirely with the basic baryon  $b$  if we assign to the triplet  $t$  the following properties: spin  $\frac{1}{2}$ ,  $z = -\frac{1}{3}$ , and baryon number  $\frac{1}{3}$ . We then refer to the members  $u\frac{2}{3}$ ,  $d-\frac{1}{3}$ , and  $s-\frac{1}{3}$  of the triplet as "quarks"  $q$  and the members of the anti-triplet as anti-quarks  $\bar{q}$ . Baryons can now be constructed from quarks by using the combinations  $(qqq)$ ,  $(qqq\bar{q})$ , etc., while mesons are made out of  $(q\bar{q})$ ,  $(q\bar{q}\bar{q})$ , etc. It is assuming that the lowest baryon configuration  $(qqq)$  gives just the representations **1**, **8**, and **10** that have been observed, while the lowest meson configuration  $(q\bar{q})$  similarly gives just **1** and **8**.

6) James Joyce, Finnegans Wake: „three quarks for Muster Mark“

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QCD is a non-Abelian gauge theory with gauge group  $G = SU(3)_c$  ( $c$  for *color*)
- 4 Matter fields of QCD (quarks) are fermions with spin 1/2, which show up in six different *flavors*

## Light quarks

<i>flavor</i>	<i>u</i>	<i>d</i>	<i>s</i>
mass [MeV]	$2.2^{+0.6}_{-0.4}$	$4.7^{+0.5}_{-0.4}$	$96^{+8}_{-4}$
charge [ $e > 0$ ]	$\frac{2}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$
$I_3$	$+\frac{1}{2}$	$-\frac{1}{2}$	0
			strangeness: $-1$

## Heavy quarks

<i>flavor</i>	<i>c</i>	<i>b</i>	<i>t</i>
mass [GeV]	$1.28 \pm 0.03$	$4.18^{+0.04}_{-0.03}$	$173.1 \pm 0.6$
charge [ $e > 0$ ]	$\frac{2}{3}$	$-\frac{1}{3}$	$\frac{2}{3}$
$I_3$	0	0	0
	charm: +1	bottom: -1	top: +1

See <http://pdg.lbl.gov>

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- General  $N_c$ :

$$\frac{1}{\sqrt{N_c!}} \epsilon_{i_1 \dots i_{N_c}} \chi^{i_1} \otimes \dots \otimes \chi^{i_{N_c}}$$

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- 11 Mesons:  $q\bar{q}$  (quark-antiquark) states; color neutral via  $\frac{1}{\sqrt{3}}(r\bar{r} + g\bar{g} + b\bar{b})$